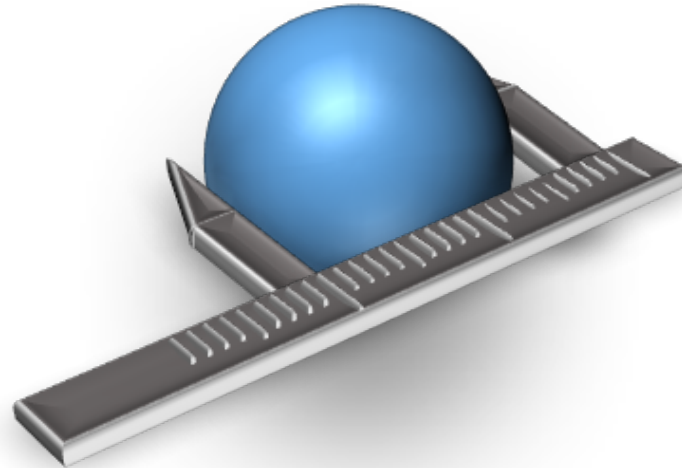


# **Proton Charge Radius**

## **by electron scattering under the lowest-ever momentum transfer**

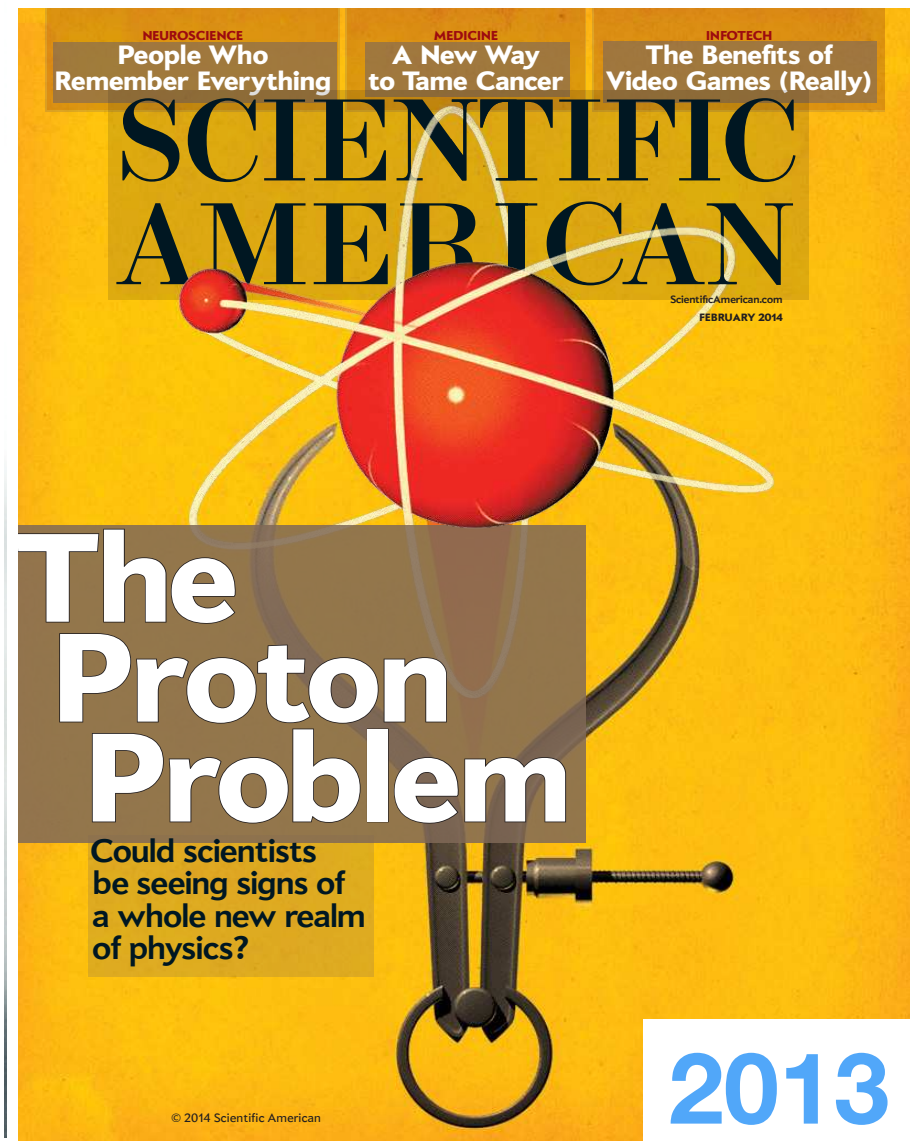


**Toshimi Suda**  
Research Center for Electron-Photon Science (ELPH),  
Tohoku University, Sendai, JAPAN

# **Proton Charge Radius Puzzle**



R. Pohl *et al.*,  
Nature 466 (2010) 213.



A. Antognini *et al.*,  
Science 339 (2013) 417.

***many many discussions ..***

electron  
scattering

***Data ? Interpretation ?***

~)

***Higher order effects ?***

hydrogen  
spectroscopy  
(990~)

***QED calculation ?***

0.82 ***New Physics (beyond SM ?)***

Proton Charge Radius (fm)

***not yet settled***

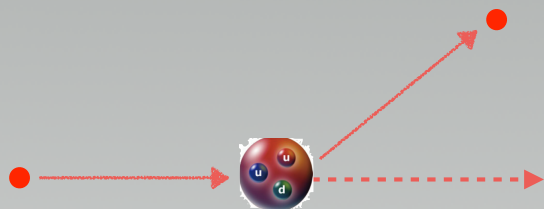
***New experiments ...***



# Proton Charge Radius measurements

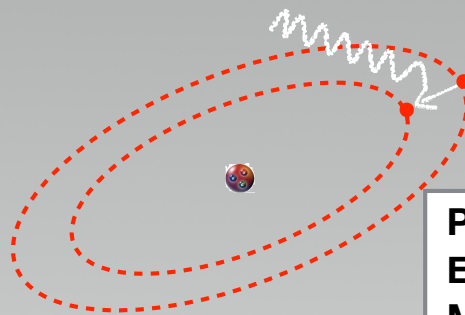
FPUA @ Nagoya  
Jan. 8-9, 2018

Electron Scattering (1950~)



$\rho(r)$  or  $\langle r^2 \rangle$

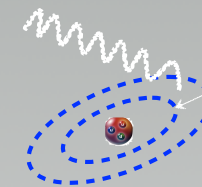
eH-Spectroscopy (1990 ~)



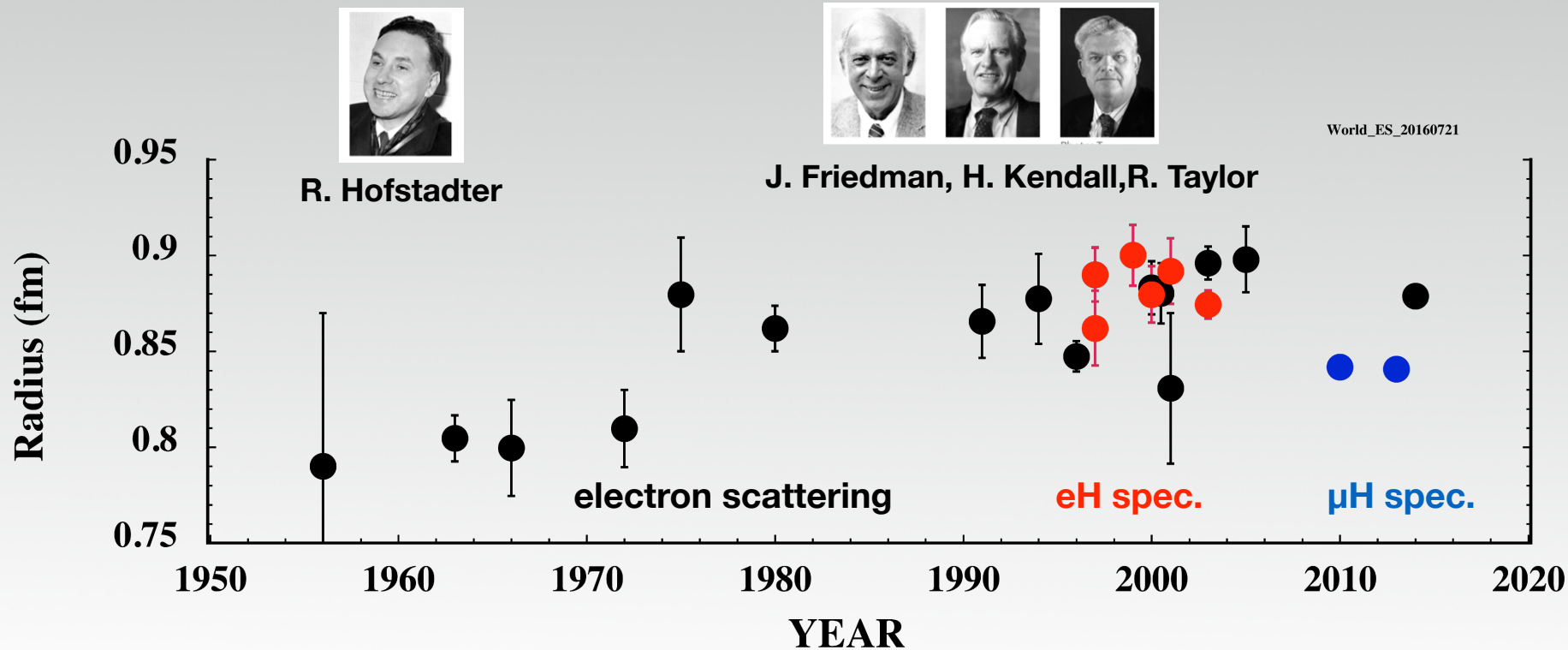
Proton Radius  $\sim 10^{-15}$  m  
Electron Orbit  $\sim 10^{-10}$  m  
Muon Orbit  $\sim 10^{-12}$  m

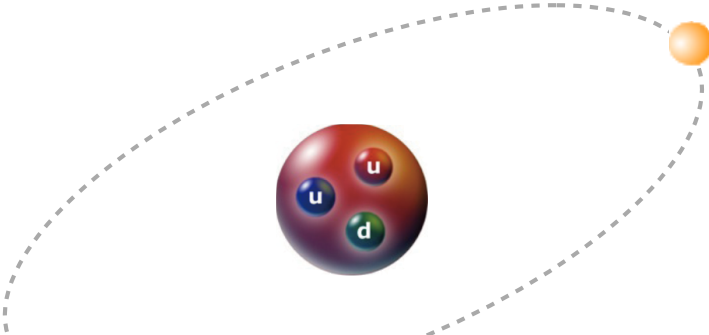
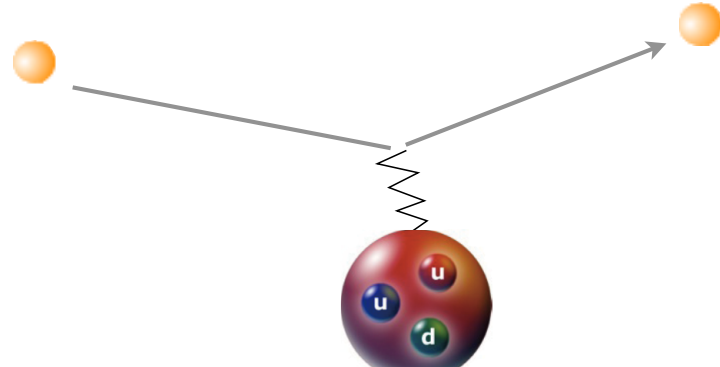
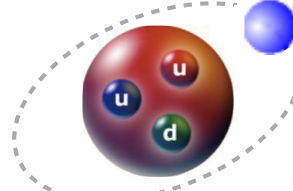
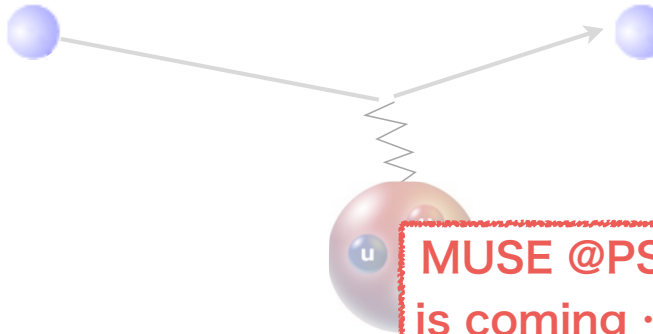
$\mu$ H-Spectroscopy (2000~)

$m_\mu \sim 200 m_e$



$$\Delta E = \alpha \cdot R_{Rydberg} + \beta \cdot \langle r^2 \rangle$$



	Spectroscopy	Scattering
$e^-$	 <p>0.8758(77)</p>	 <p>0.8770(60)</p>
$\mu^-$	 <p>0.8409(4)</p>	 <p>MUSE @PSI is coming ...</p>

$m_e = 0.511 \text{ MeV}$   
 $m_\mu = 105.6 \text{ MeV}$

**e-scattering** : re-analysis of world data

*model dependence of radius determination  
radiative correction, higher order effects ...*

**H-spec.** : QED calculations by several groups

*higher order effects such as two-photon exchange ...*



**No serious problem yet identified**

**possible non identical nature of e- $\mu$  ?**

*Lepton universality*

## New experiments

**e-scatt.**

JLAB (US) 、 Mainz (Germany) 、 **Tohoku (Japan)**

**H-spec.**

Toronto (Canada) 、 Garching (Germany)  
LKB (France) 、 NIST (US) etc.

**$\mu$ H-spec**

PSI (Switzerland)、 RIKEN (Zemach moment)

**$\mu^\pm$ scatt.**

PSI (Switzerland)

# Electron scattering group of Tohoku Univ.

World's first

## 1) short-lived exotic nuclei @ RIKEN-SCRIT facility

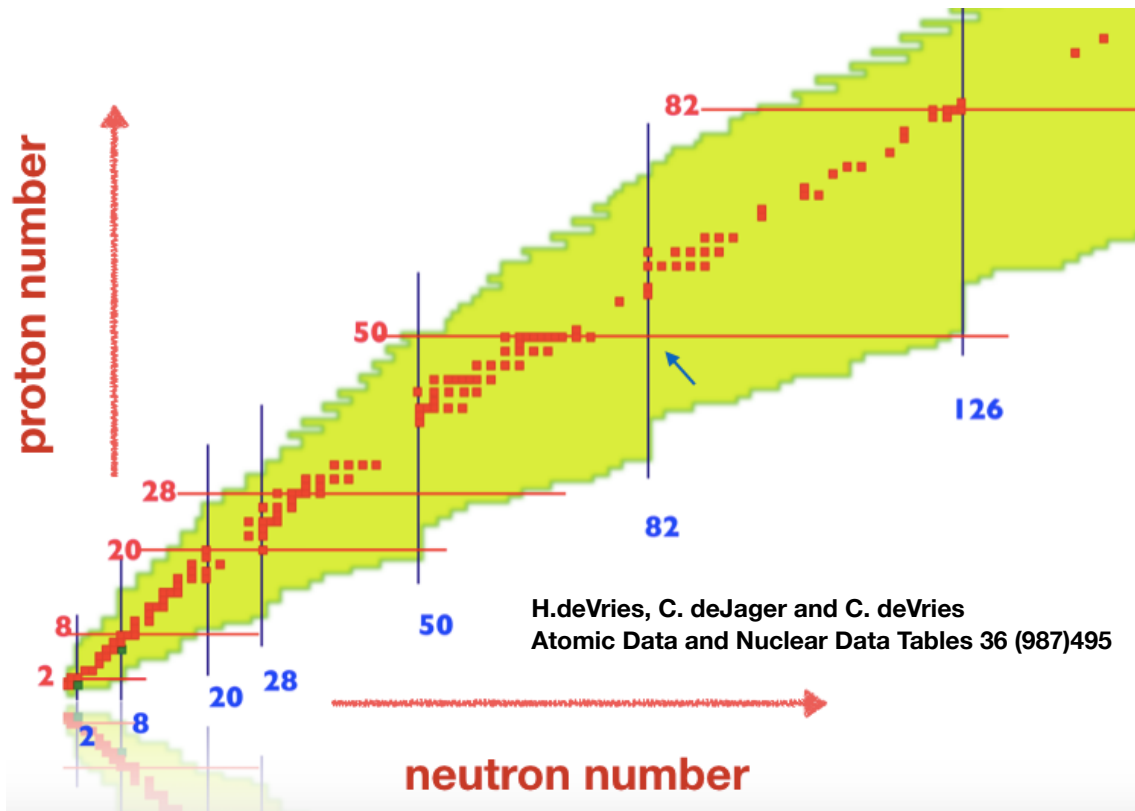
World's first electron scattering facility

*charge radii and charge density distributions*

## 2) proton @ Tohoku Univ. (Sendai)

World's lowest  $Q^2$

## Nuclei targeted for electron scattering



## Short-lived Exotic Nuclei

Production-hard + Short-lived

## Elastic electron scattering

$$\frac{d\sigma}{d\Omega} = \frac{d\sigma_{\text{Mott}}}{d\Omega} |F_c(q)|^2$$

$$F_c(q) = \int \rho_c(\vec{r}) e^{i\vec{q}\vec{r}} d\vec{r}$$

$$\rho_c(\vec{r}) = \sum_p \psi_p^*(\vec{r}) \psi_p(\vec{r})$$

## SCRIT ( Self-Confining RI ion Target )

$L \sim 10^{27} / \text{cm}^2/\text{s}$  with only  $\sim 10^8$  target nuclei

target thickness used to be  $\sim 10^{20} / \text{cm}^2/\text{s}$

## Low Luminosity

Charge density distribution  
Charge radius

## SCRIT collaboration

RIKEN + Rikkyo U. + Tohoku U.

(Review) T. Suda and H. Simon, *Prog. Part. Nucl. Phys.* 96 (2017) 1-31.

(First Physics Data) K. Tsukada et al., *Phys. Rev. Lett.* 118(2017) 262501.



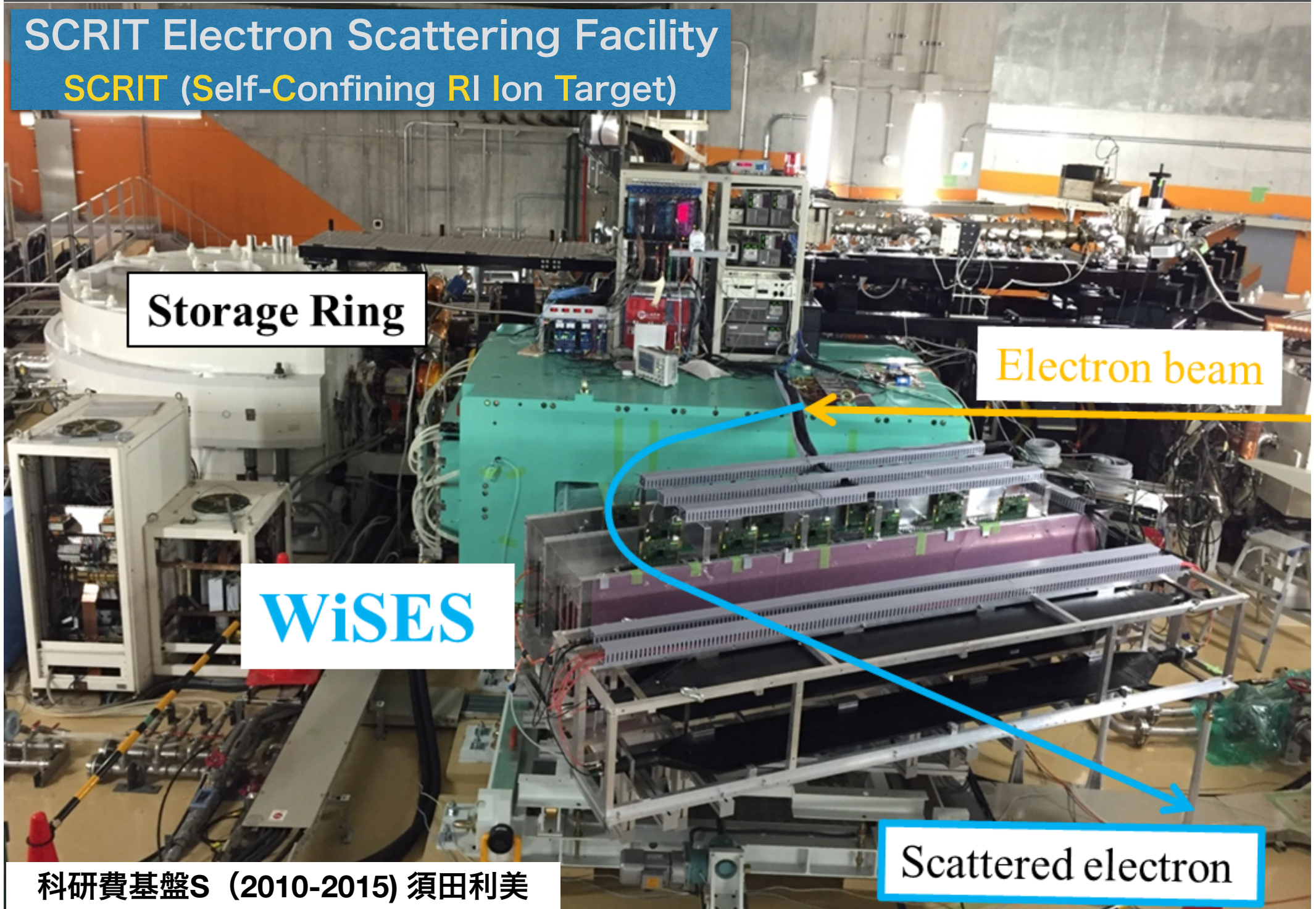
## SCRIT Electron Scattering Facility SCRIT (Self-Confining RI Ion Target)

Storage Ring

Electron beam

WiSES

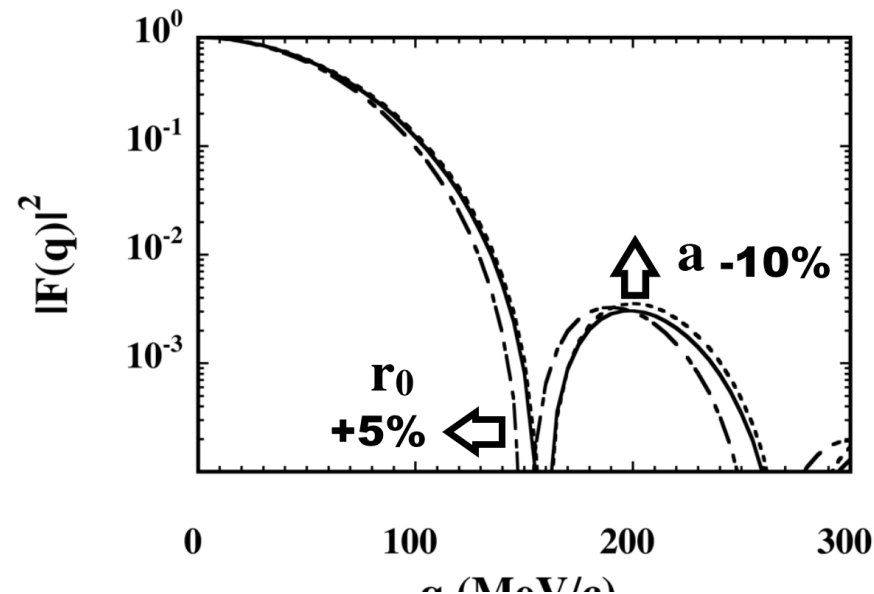
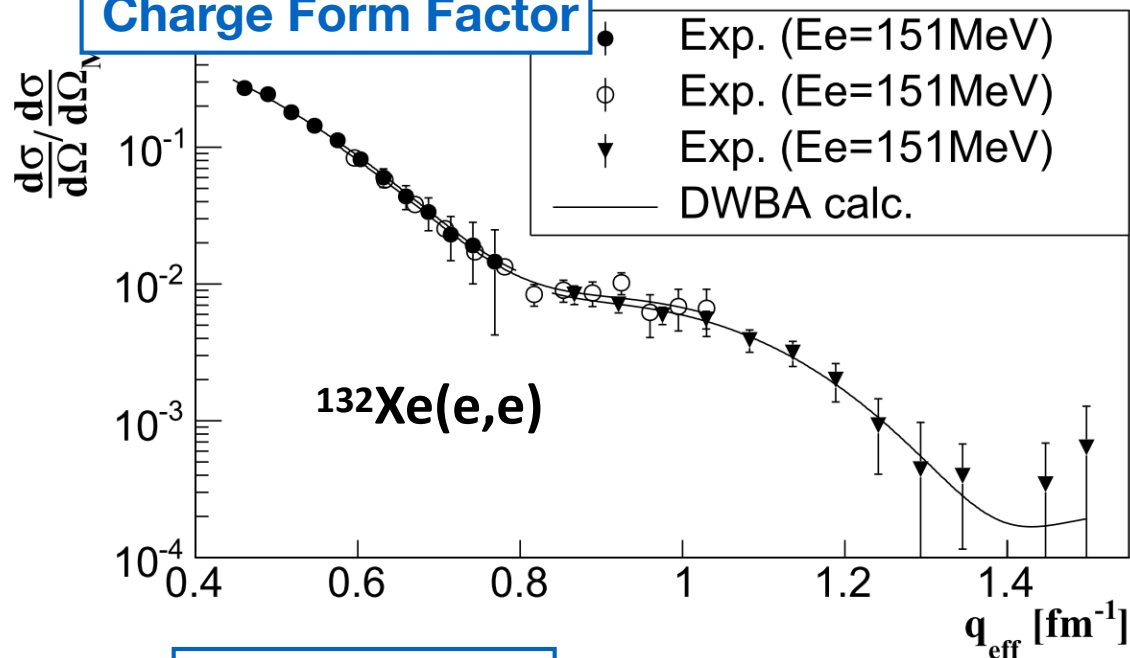
Scattered electron



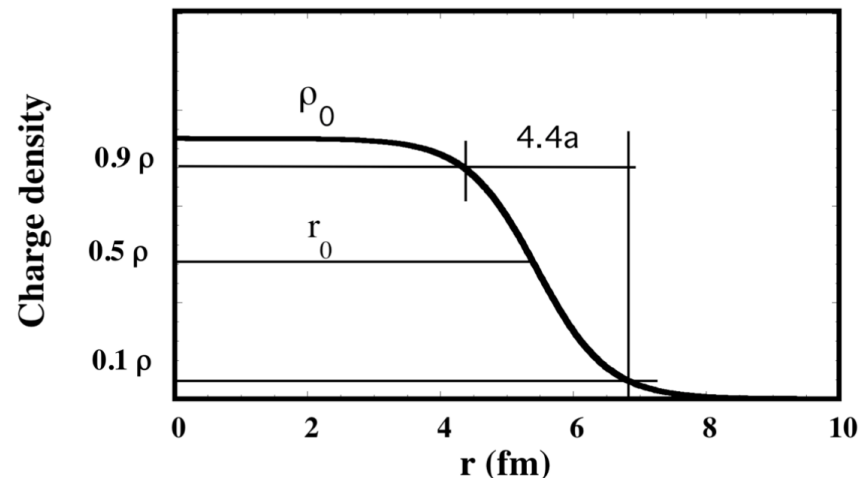
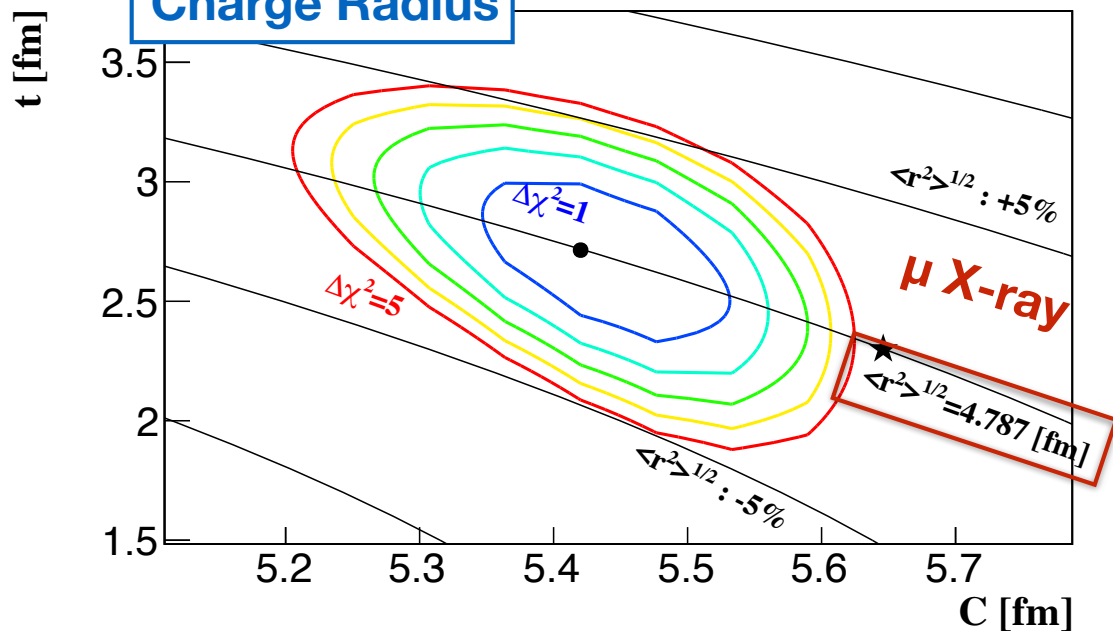


using only  $\sim 10^8$  target nuclei

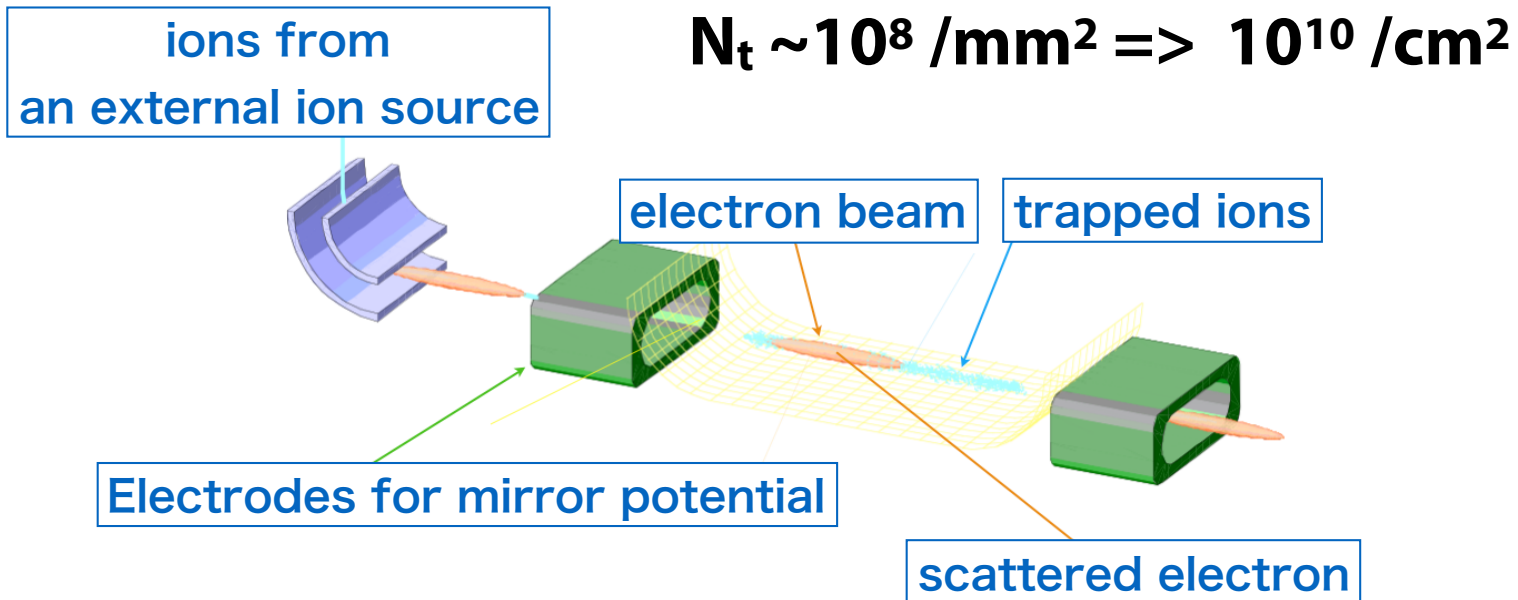
## Charge Form Factor



## Charge Radius



$\sim 10^8$  ions are trapped on e-beam ( $\sim 1 \text{ mm}^2$ )



	<b>E<sub>e</sub></b>	<b>N<sub>beam</sub></b>	<b><math>\rho \cdot t</math></b>	<b>L</b>
<b>Hofstadter's era (1950s)</b>	150 MeV	$\sim 1 \text{ nA}$ ( $\sim 10^9 / \text{s}$ )	$\sim 10^{19} / \text{cm}^2$	$\sim 10^{28} / \text{cm}^2 / \text{s}$
<b>JLAB</b>	6 GeV	$\sim 100 \mu\text{A}$ ( $\sim 10^{14} / \text{s}$ )	$\sim 10^{22} / \text{cm}^2$	$\sim 10^{36} / \text{cm}^2 / \text{s}$
<b>SCRIT</b>	<b>150 - 300 MeV</b>	<b><math>\sim 200 \text{ mA}</math></b> <b>(<math>\sim 10^{18} / \text{s}</math>)</b>	<b><math>\sim 10^{10} / \text{cm}^2</math></b>	<b><math>\sim 10^{27} / \text{cm}^2 / \text{s}</math></b>

The SCRIT technique reduces the required target thickness by  $\sim 10^{-10}$  !!

# **Proton Charge Radius**

## RMS radius

$$\begin{aligned}\langle r^2 \rangle &= \int r^2 \rho(\vec{r}) d\vec{r} \\ &= 4\pi \int r^4 \rho(r) dr\end{aligned}$$

## $\rho(r)$

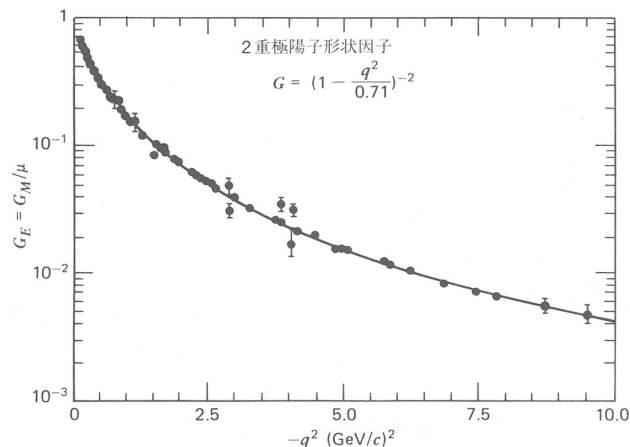
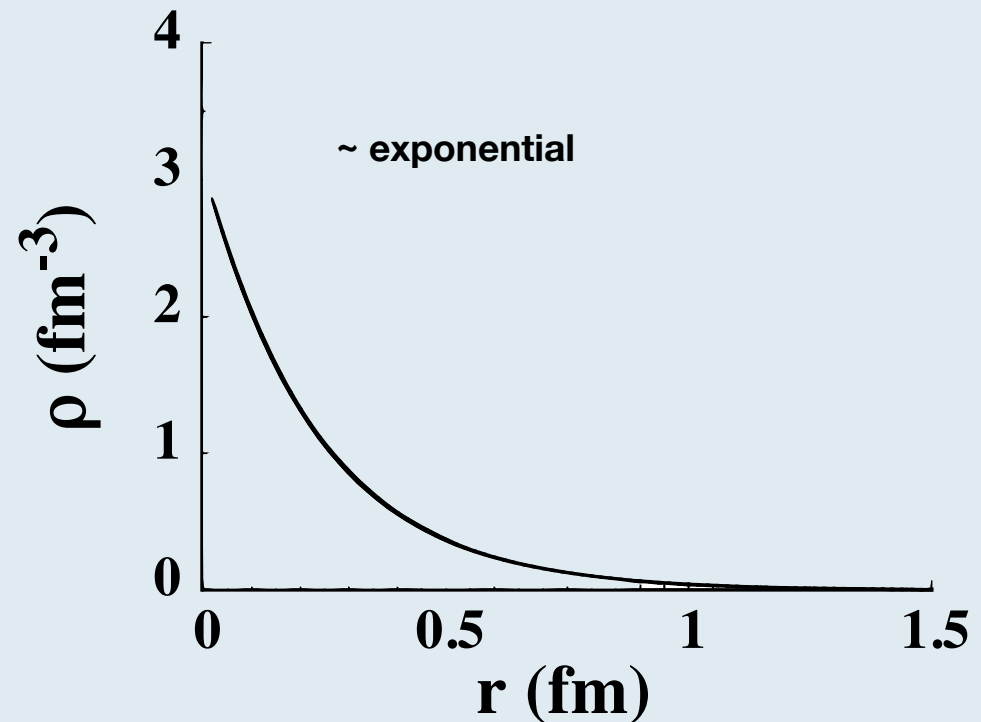
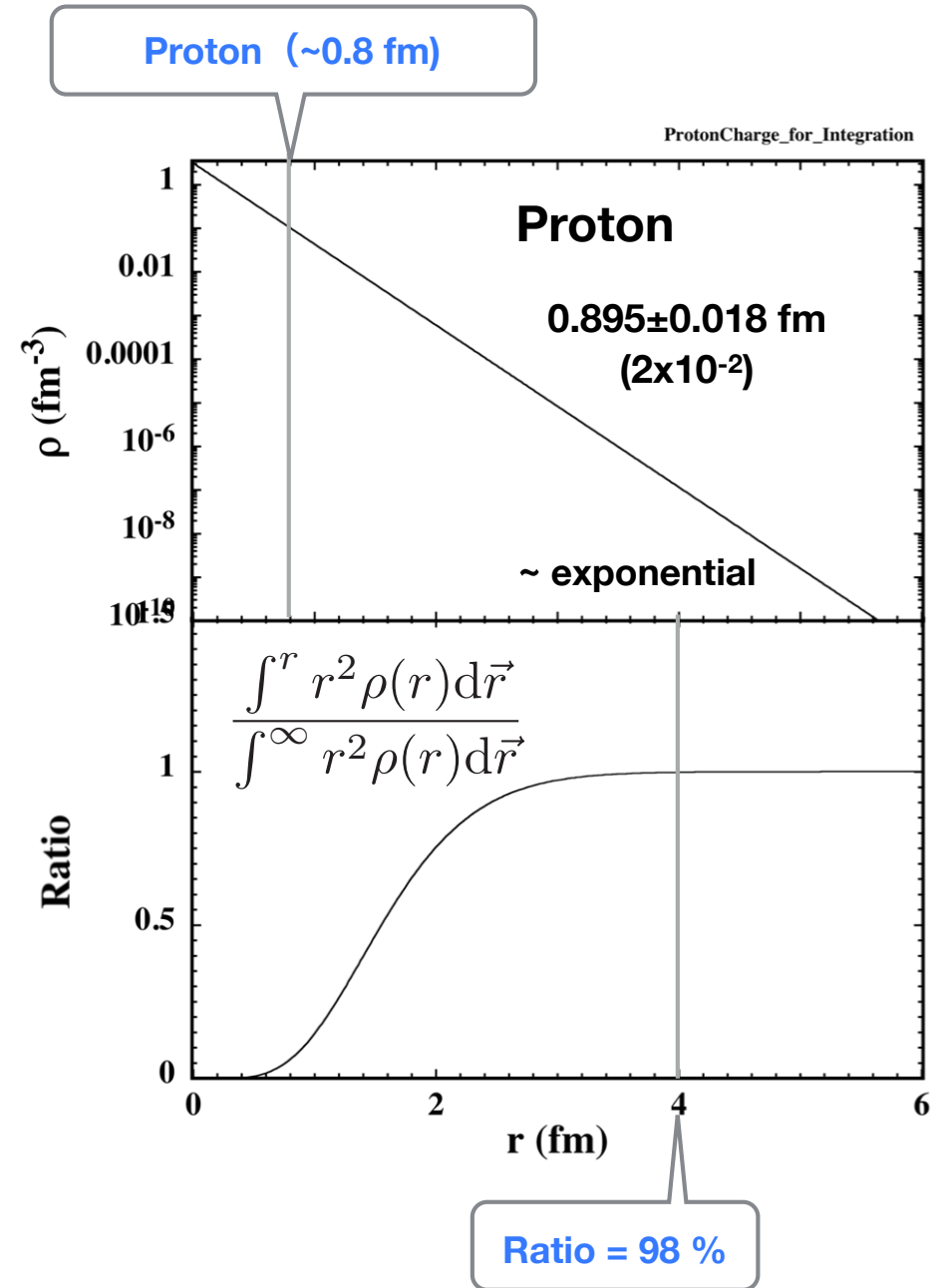
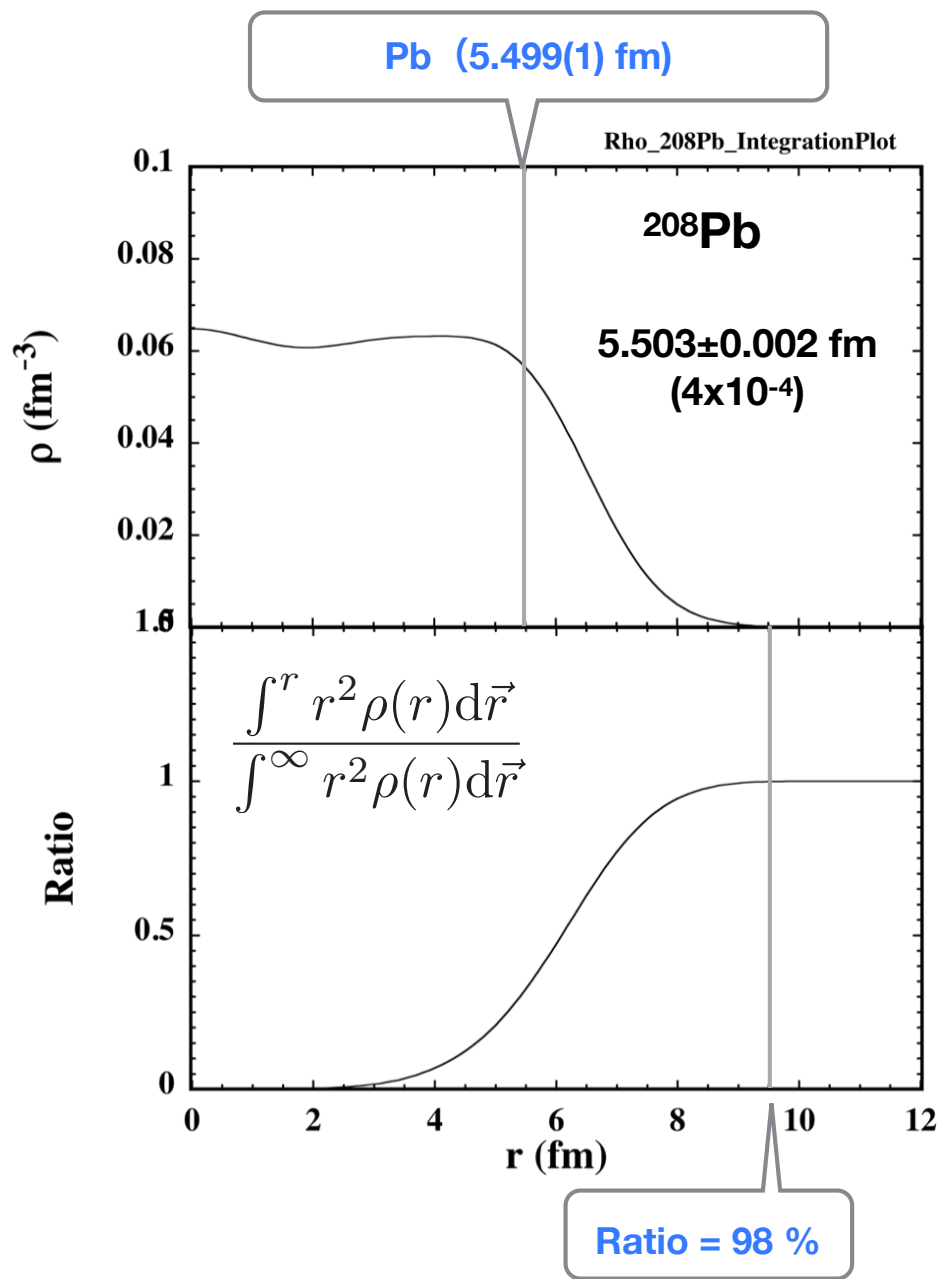


図 8・4  $q^2$  の関数としての陽子形状因子



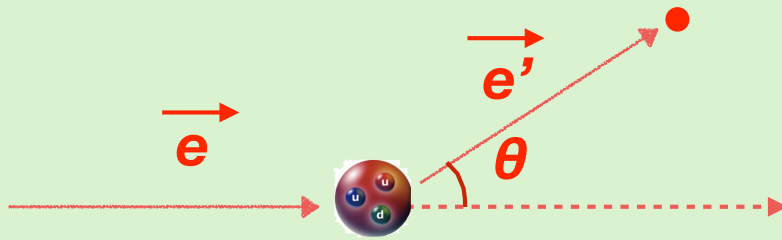
# Charge Radius and Density Distribution

**Charge Radius**  $\langle r^2 \rangle = \int r^2 \rho(\vec{r}) d\vec{r}$



**Charge Radius**  
**by**  
**Electron Scattering**





$$\frac{d\sigma}{d\Omega} = \left(\frac{d\sigma}{d\Omega}\right)_{Mott} \frac{G_E^2(Q^2) + \frac{\tau}{\epsilon} G_M^2(Q^2)}{1 + \tau}$$

$$\left(\frac{d\sigma}{d\Omega}\right)_{Mott} = \frac{z^2 \alpha^2 \cos^2(\theta/2)}{4e^2 \sin^4(\theta/2)} \propto \frac{e^2}{q^4}$$

momentum transfer  $\vec{q} = \vec{e} - \vec{e}'$

energy transfer  $\omega = e - e'$

4 momentum transfer  $Q^2 = q^2 - \omega^2$   
 $= 4 e e' \sin^2(\theta/2)$

$$\epsilon = \frac{1}{1 + 2(1 + \tau)\tan^2\frac{\theta}{2}}$$

$$\tau = \frac{Q^2}{4m_p^2}$$

**Under One-Photon Exchange approximation**

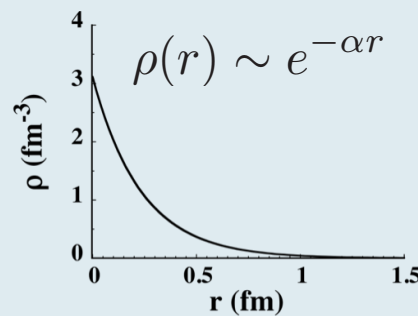
## 1) high $Q^2$ : charge density $\rho(r)$

Electric Form Factor  $G_E$



$\rho(r)$

$$\langle r^2 \rangle = \int r^2 \rho(\vec{r}) d\vec{r}$$



radius is sensitive to  $\rho(r)$  at large distance  
( even at  $r \sim 4$  fm )

## 2) low $Q^2$

$$G_E(Q^2) \sim 1 - \frac{\langle r^2 \rangle^{1/2}}{6} Q^2 + \frac{\langle r^4 \rangle^{1/2}}{120} Q^4 - \dots$$

$$\langle r^2 \rangle \equiv -6 \left. \frac{dG_E(Q^2)}{dQ^2} \right|_{Q^2 \rightarrow 0}$$

**ill problem : higher order contribution**



**lower  $Q^2$  as possible**

## Elastic scattering off proton

$$\frac{d\sigma}{d\Omega} \propto G_E^2(Q^2) + \alpha(\theta) G_M^2(Q^2)$$

Charge F.F.

Magnetic F.F.

### 1 Rosenbluth separation for $G_E(Q^2)$

measurements at various  $\theta$  under fixed  $Q^2$

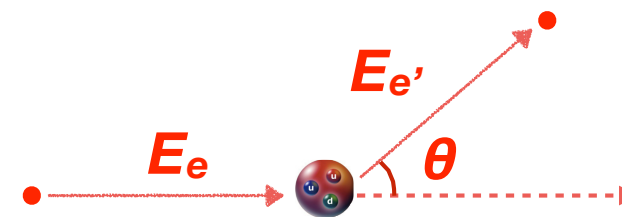
➔ frequent change of  $E_e$  required

➔  $G_E(Q^2)$  determination for various  $Q^2$

### 2 $G_E(Q^2)$ and Charge Radius

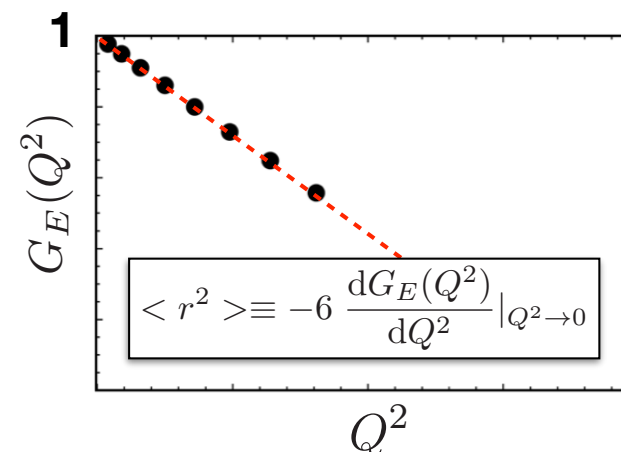
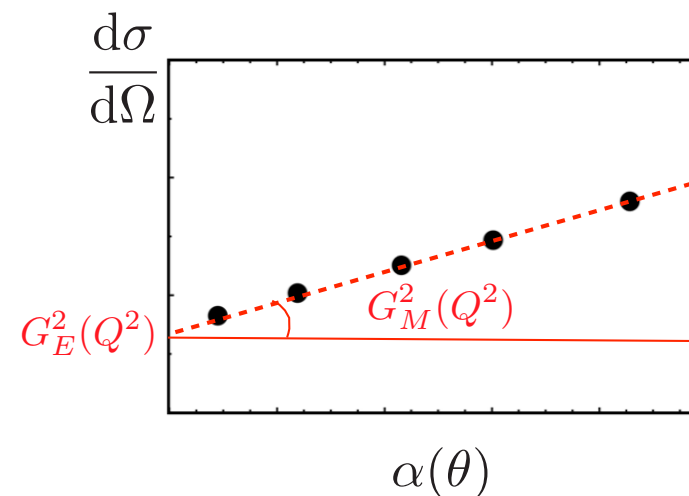
$$G_E(Q^2) \sim 1 - \frac{\langle r^2 \rangle}{6} Q^2 + \frac{\langle r^4 \rangle}{120} Q^4 - \frac{\langle r^6 \rangle}{5040} Q^6 + \dots$$

Taylor expansion of  $G_E(Q^2)$  at low  $Q^2$



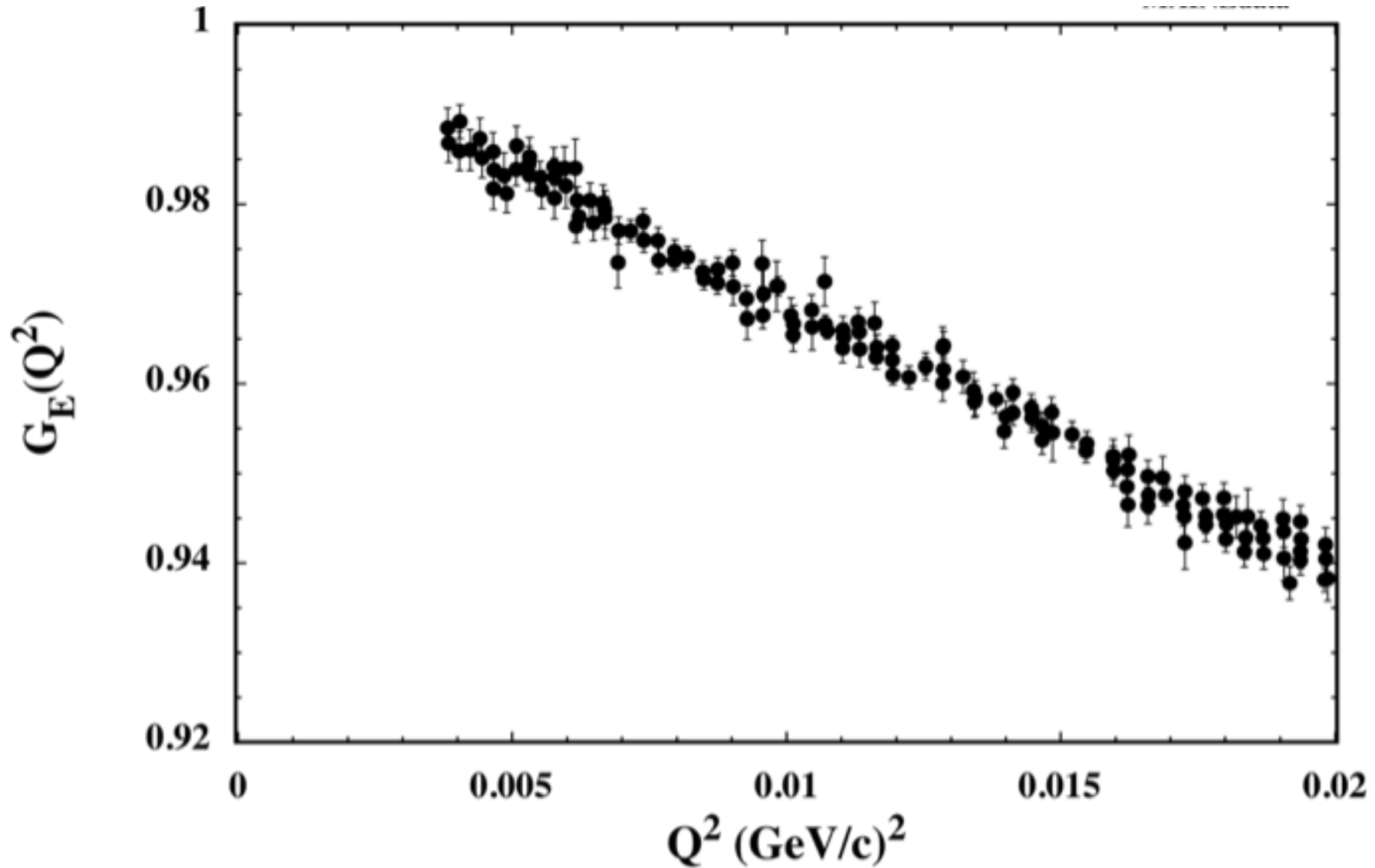
4 momentum transfer

$$Q^2 = 4 E_e E_e' \sin^2(\theta/2)$$

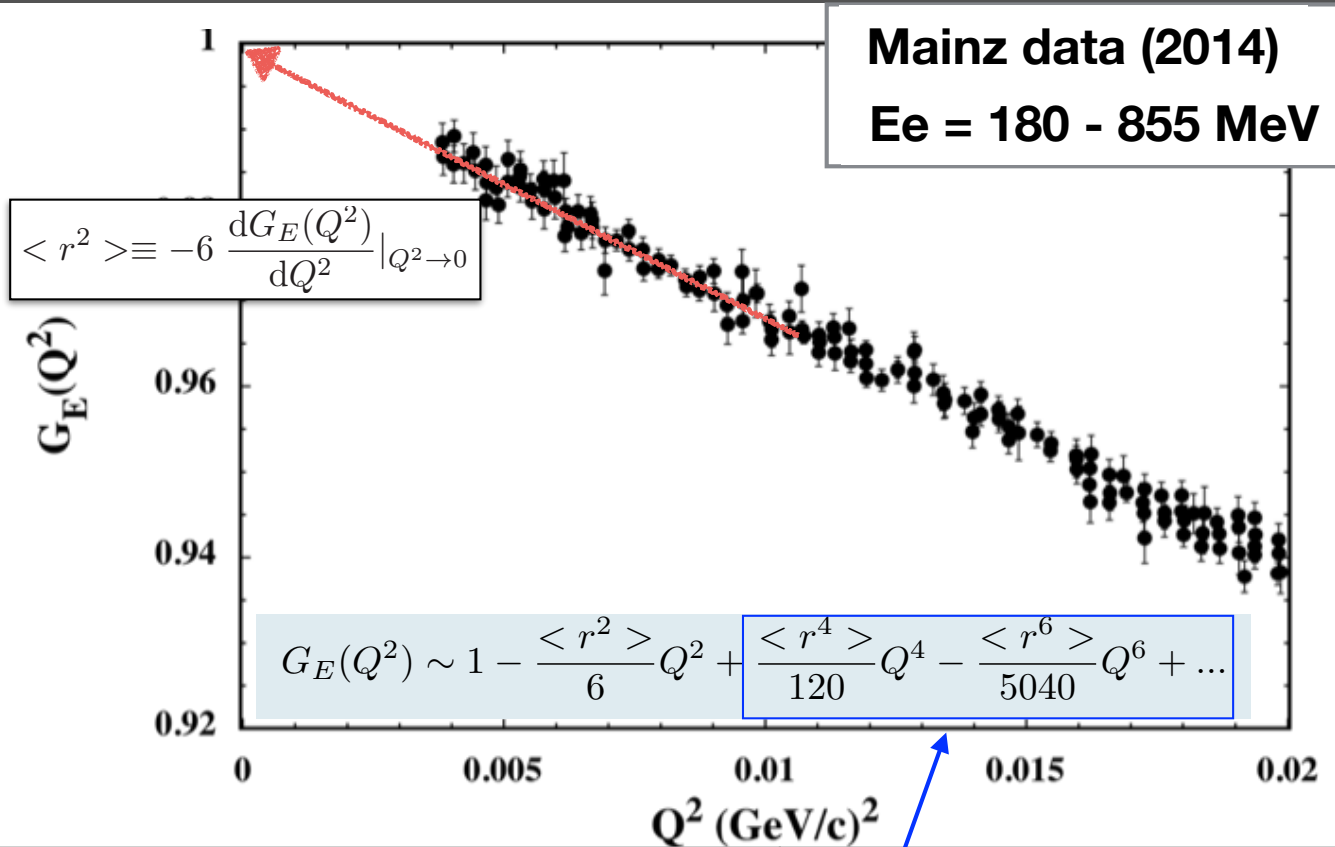


## Mainz data

$E_e = 180 - 855$  MeV



# Deduced Radius from Mainz data



potential problems	backgrounds
No very low- $Q^2$ data	high-energy electron accelerator
No Rosenbluth separation	frequent $E_e$ change : impractical
No absolute cross section	spatially extended Liq-H <sub>2</sub> target + spectrometer

treatments of higher-order effects



introduce large uncertainty in the radius

(Phys. Rev. C90 (2014) 045206)

# **Electron Scattering at Lowest-ever Momentum Transfer**

## **ULQ<sup>2</sup> (Ultra-Low Q<sup>2</sup>) Collaboration**

Tohoku Univ. : H. Kikunaga, K. Tsukada, Y. Honda, T. Tamae, T. Mutoh, K. Takahashi,

K. Nanbu, M. Miyabe, A. Tokiyasu, K. Nanba, T. Aoyagi

S. Sasaki, N. Tsukamoto

Miyazaki Univ. : Y. Maeda

Hampton Univ. : Michael Kohl

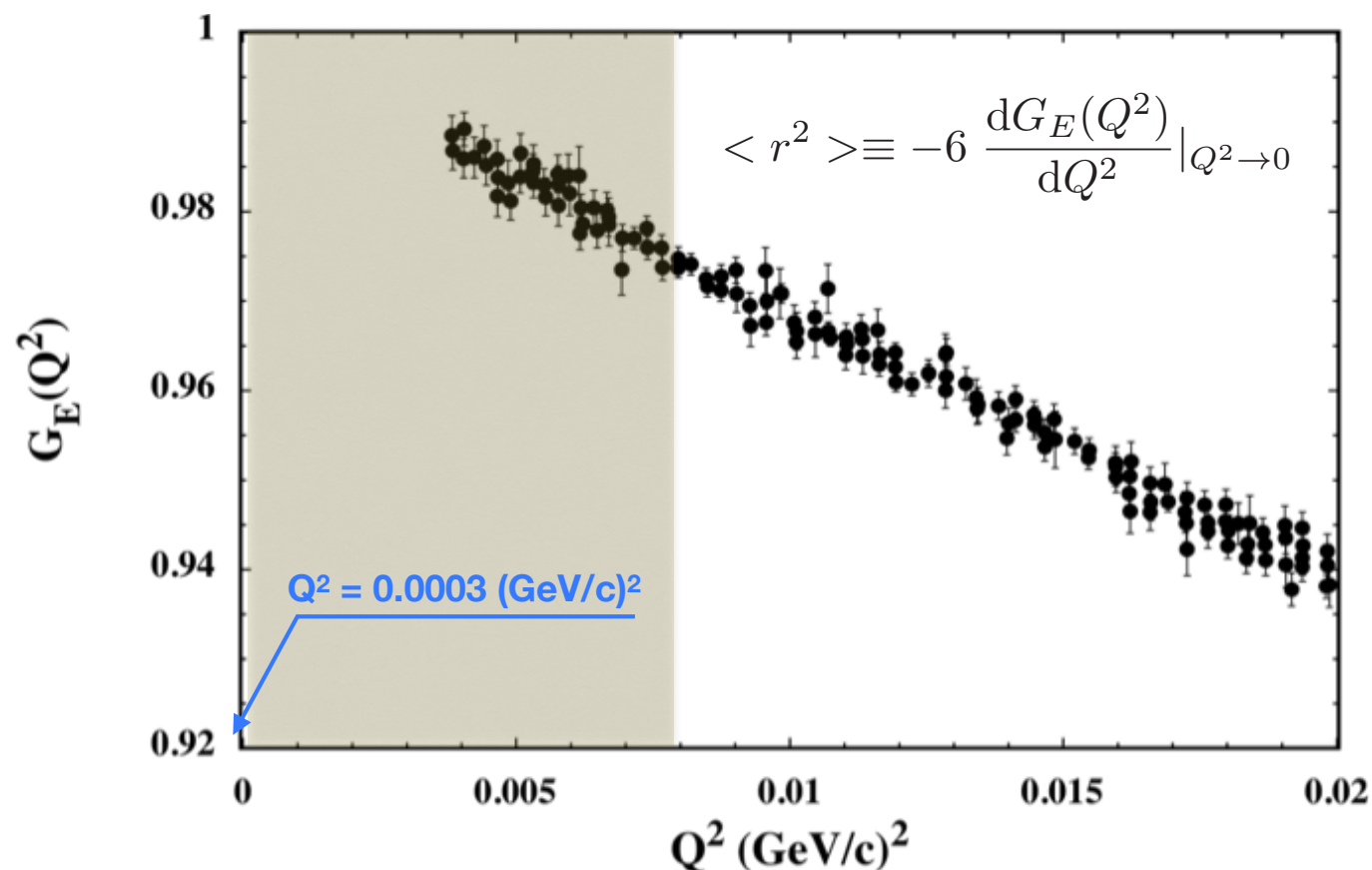
## Our Goal !

$G_E(Q^2)$  measurements at  $0.0003 \leq Q^2 \leq 0.008$  (GeV/c)<sup>2</sup>

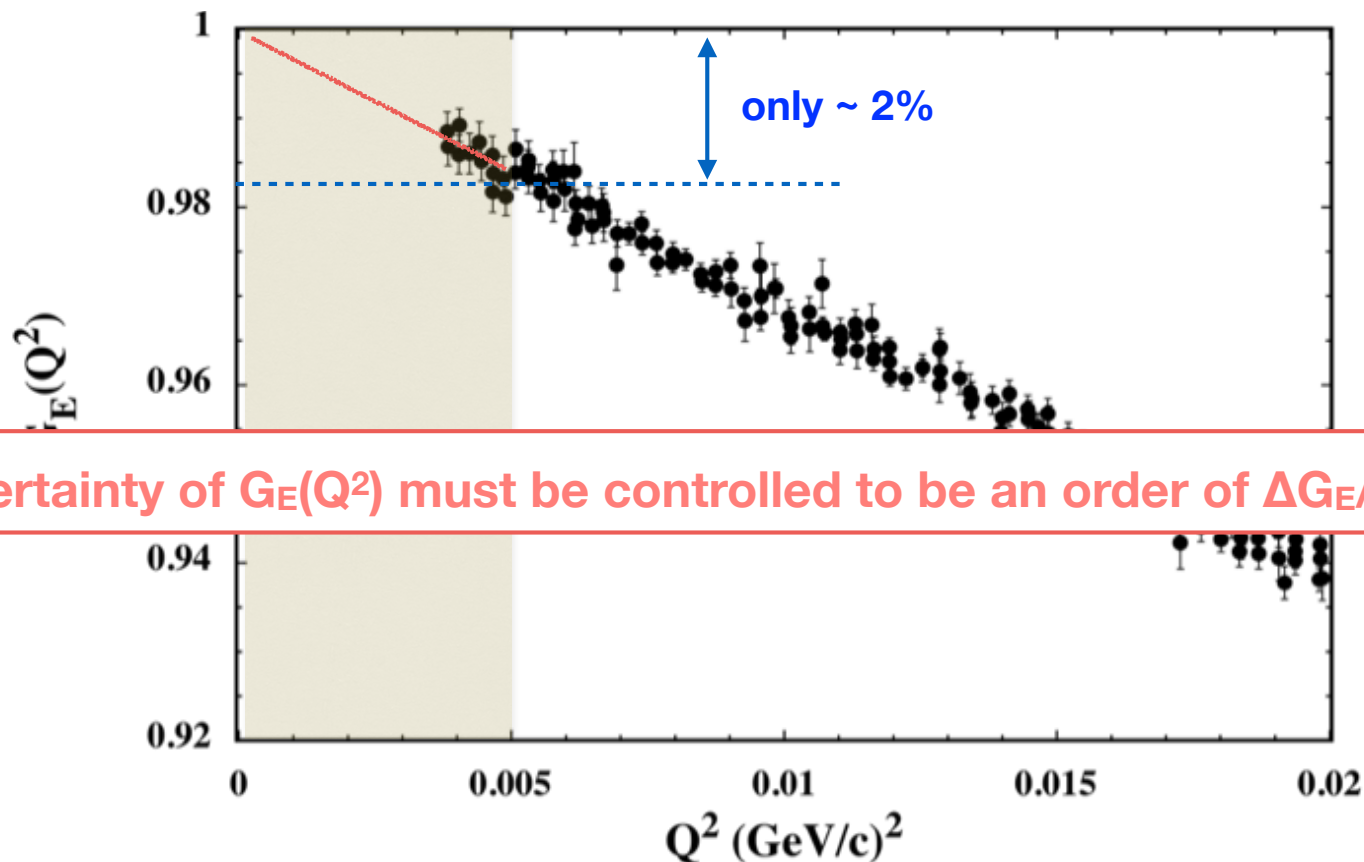
Low energy electron beam (  $20 \leq E_e \leq 60$  MeV)

**Absolute** cross section measurement

**Rosenbluth separation** ( $G_E(Q^2)$ ,  $G_M(Q^2)$  separation)







Uncertainty of  $G_E(Q^2)$  must be controlled to be an order of  $\Delta G_E/G_E \sim 10^{-3}$

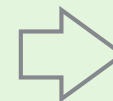
$$\frac{dN_{evt}}{d\Omega} = \frac{d\sigma}{d\Omega} \underbrace{N_{target}}_{\text{target thickness}} \underbrace{N_{beam}}_{\text{beam dose}} \underbrace{\Delta\Omega}_{\text{spectrometer acceptance}}$$

Statistics : at least  $> 10^6$  for each  $(E_e, \theta)$  measurements

Target thickness

Beam dose at various intensities

Acceptance at various scattering angle



accuracy of  $\sim 10^{-3}$   
not obvious !

## Relative measurement for $^{12}\text{C}(e,e)^{12}\text{C}$ and $p(e,e)p$

$$\frac{dN^{e^{12}\text{C}}/d\Omega}{dN^{ep}/d\Omega} = \frac{d\sigma^{e^{12}\text{C}}/d\Omega}{d\sigma^{ep}/d\Omega} \cdot \frac{N_{target}^{12\text{C}}}{N_{target}^H}$$

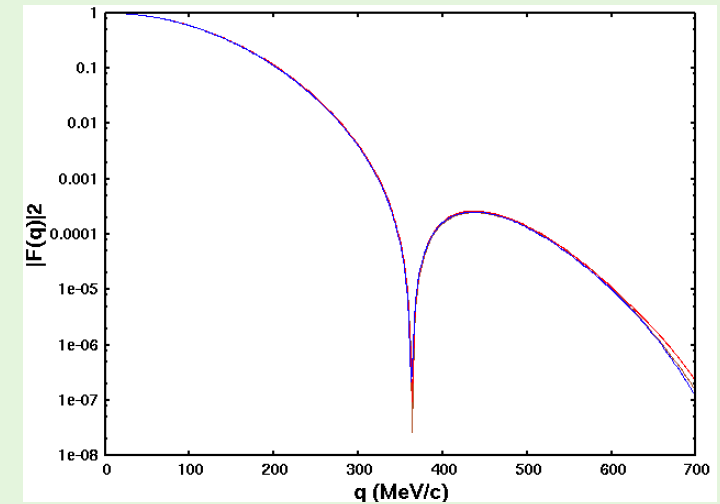
$$\frac{dN_{evt}}{d\Omega} = \frac{d\sigma}{d\Omega} N_{target} \boxed{N_{beam} \Delta\Omega}$$

Canceled out in relative measurements

## 1) <sup>12</sup>C : “standard” nucleus for (e,e')

**μ-Xray  
electron scattering**

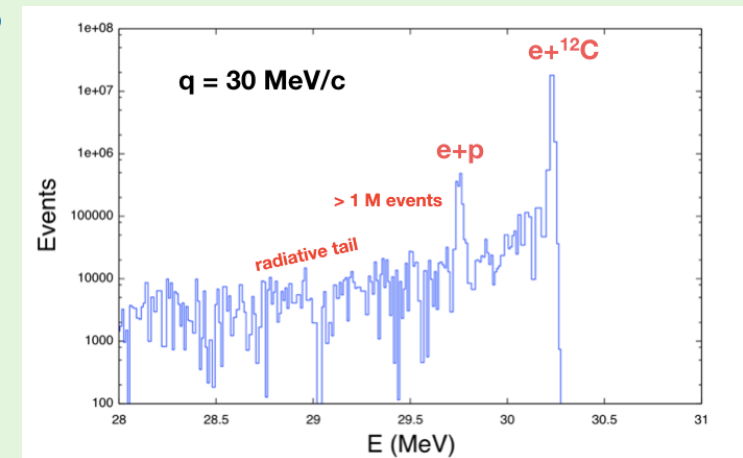
$$\frac{\Delta \langle r_{12C}^2 \rangle^{1/2}}{\langle r_{12C}^2 \rangle^{1/2}} \sim 3 \times 10^{-3}$$



## 2) <sup>12</sup>C(e,e)<sup>12</sup>C, p(e,e)p by kinematics

**ΔE = 0.2 - 4 MeV  
for q = 20 - 90 MeV/c**

**Δp/p ~ 10<sup>-3</sup>**

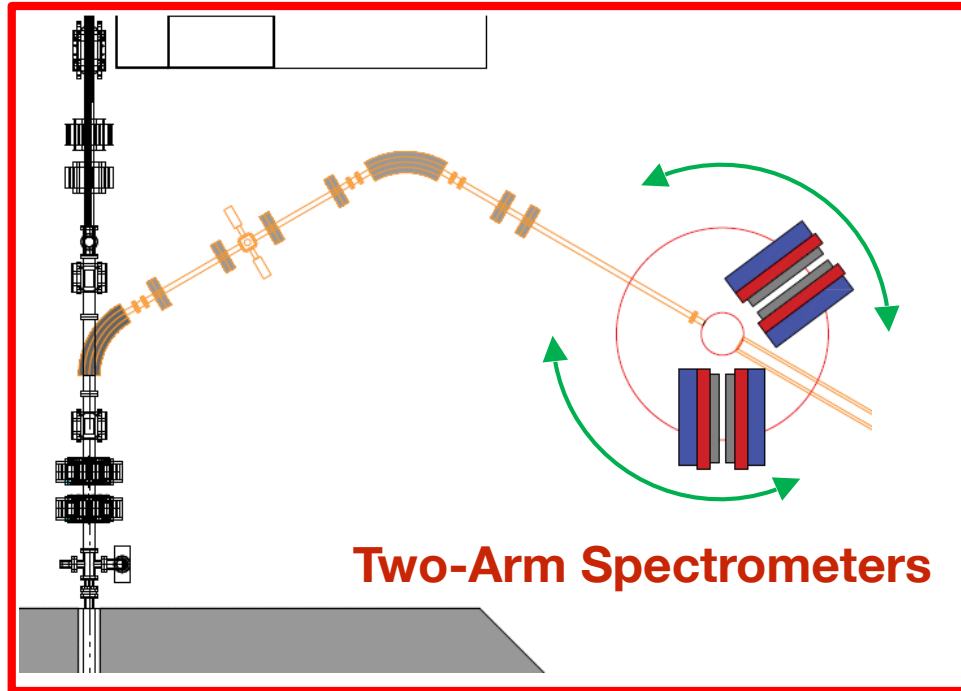


## 3) no severe damage of target is expected

large cross section :  $\frac{d\sigma}{d\Omega} \propto 1/q^4$

**I<sub>e</sub> ~ 1 nA - 1 μA**

## New beam line + Spectrometers



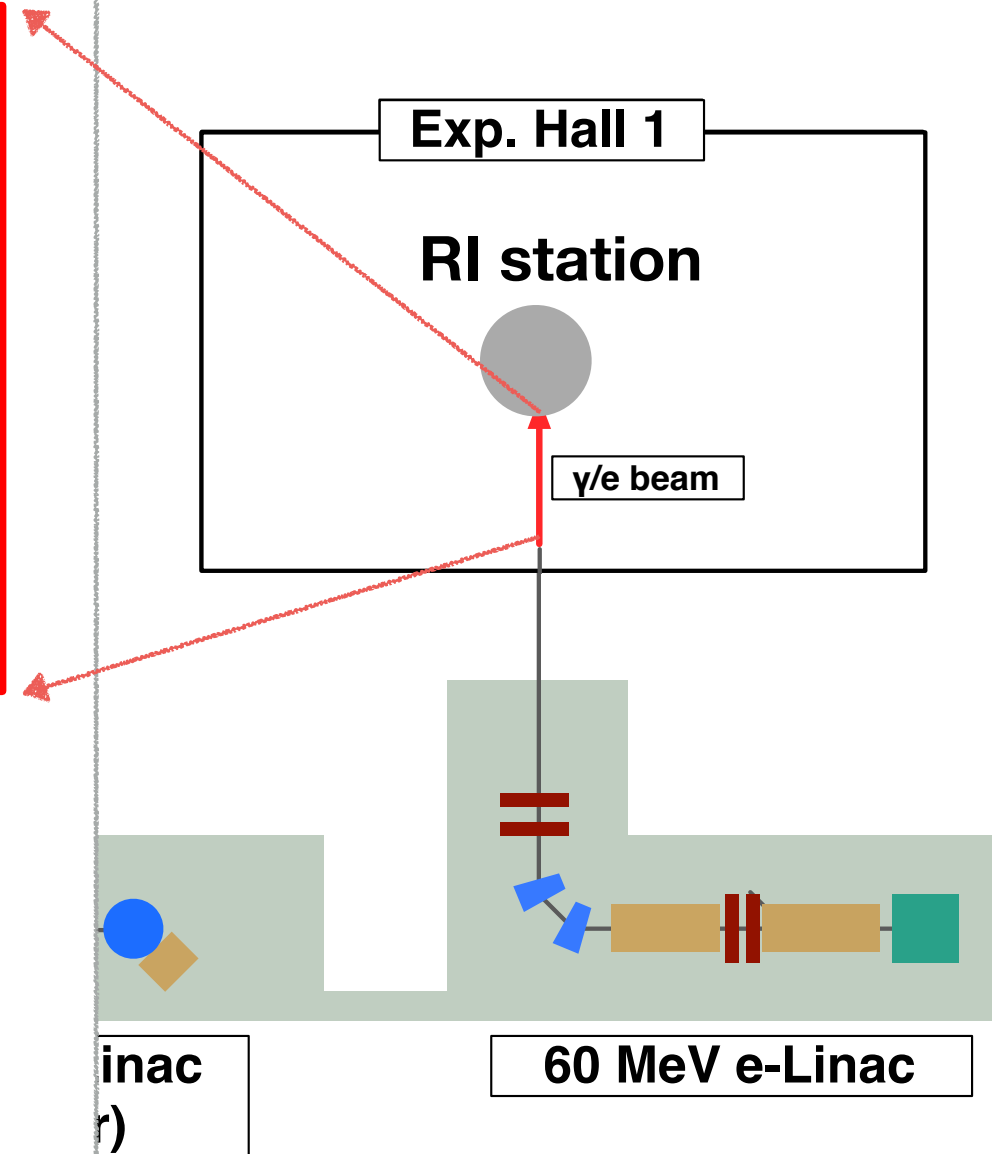
$E_e = 20 - 60 \text{ MeV}$

$I_e = 1 \text{ nA} - 0.1 \text{ } \mu\text{A}$

beam size on target  $\leq 1 \text{ mm}$

$\Delta p/p \leq 10^{-3}$

KAKENHI : Scientific Research (S) : T. Suda



## Two-Arm Electron Spectrometers

### 1) Ion-optical pro

No tracking is  
( $Pe' = 20 - 60$ )

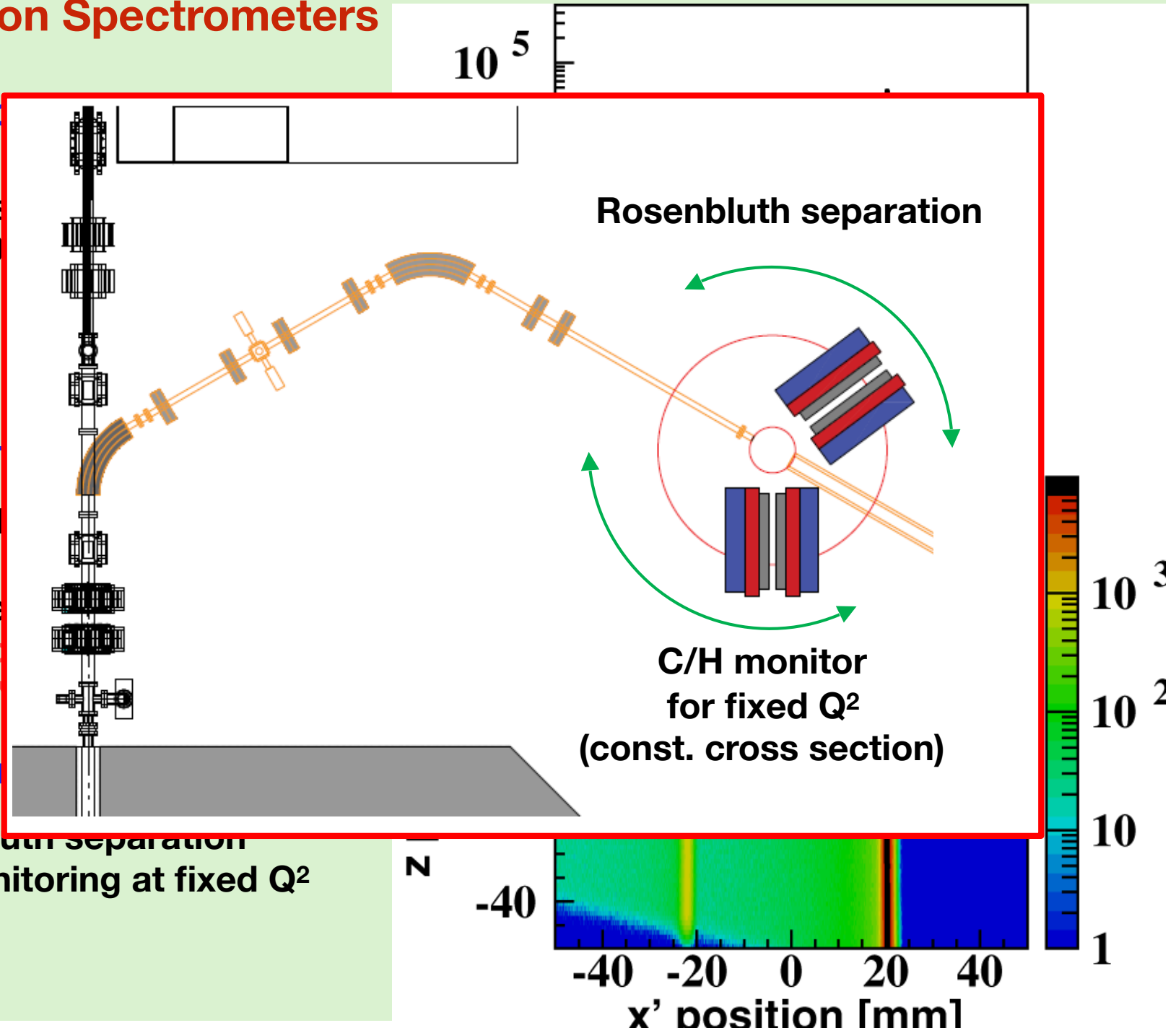
### 2) Kinematical e+

mom. res.  $\leq 1$   
mom. bite  $\sim 6$   
focal plane de

: "muc  
(T. Y

### 3) Two-arm spectr

Spec1 : Rosenbluth separation  
Spec2 : C/H monitoring at fixed  $Q^2$



Science Oct. 6, 2017

Science 06 Oct 2017:  
Vol. 358, Issue 6359, pp. 79-85  
DOI: 10.1126/science.aah6677

RESEARCH

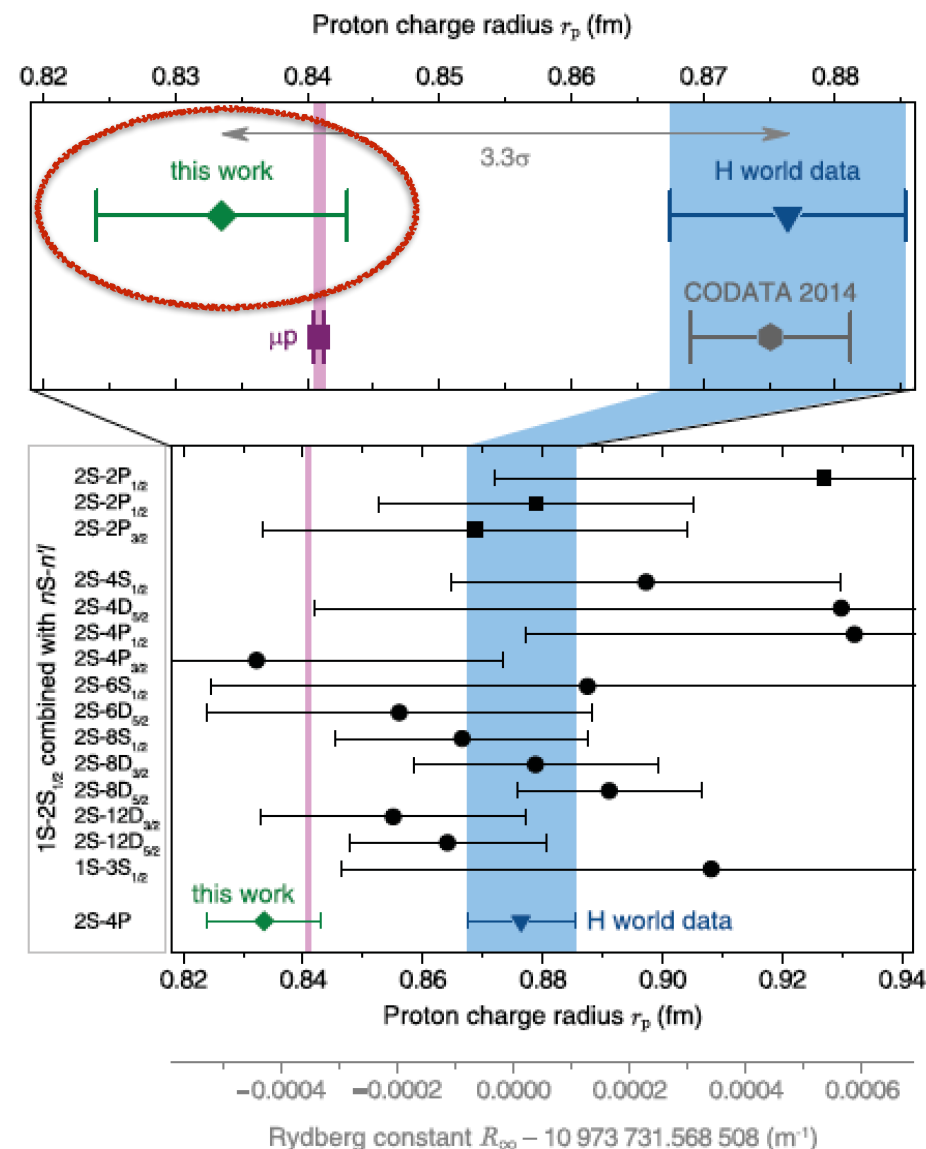
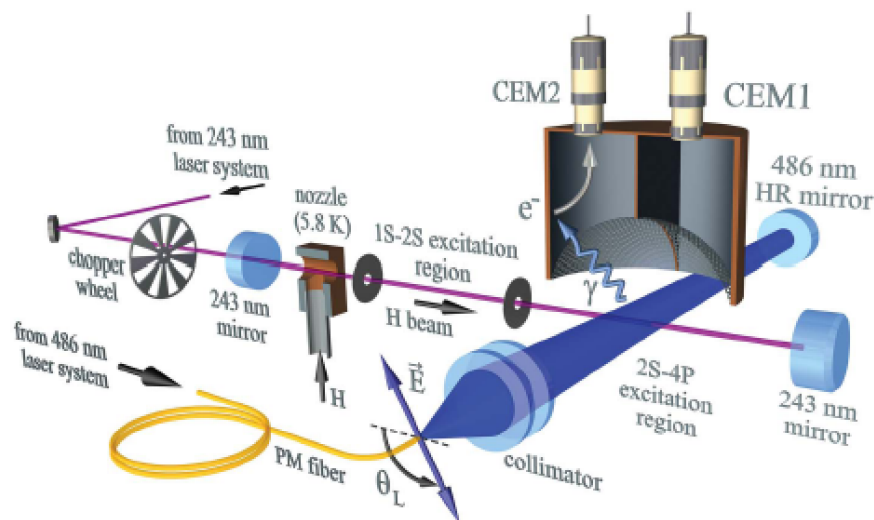
## RESEARCH ARTICLE

ATOMIC PHYSICS

# The Rydberg constant and proton size from atomic hydrogen

Axel Beyer,<sup>1</sup> Lothar Maisenbacher,<sup>1\*</sup> Arthur Matveev,<sup>1</sup> Randolph Pohl,<sup>1,†</sup>  
Ksenia Khabarova,<sup>2,3</sup> Alexey Grinin,<sup>1</sup> Tobias Lamour,<sup>1</sup> Dylan C. Yost,<sup>1,‡</sup>  
Theodor W. Hänsch,<sup>1,4</sup> Nikolai Kolachevsky,<sup>2,3</sup> Thomas Udem<sup>1,4</sup>

At the core of the "proton radius puzzle" is a four-standard deviation discrepancy between the proton root-mean-square charge radii ( $r_p$ ) determined from the regular hydrogen (H) and the muonic hydrogen ( $\mu\text{p}$ ) atoms. Using a cryogenic beam of H atoms, we measured the 2S-4P transition frequency in H, yielding the values of the Rydberg constant  $R_\infty = 10973731.568076(96)$  per meter and  $r_p = 0.8335(95)$  femtometer. Our  $r_p$  value is 3.3 combined standard deviations smaller than the previous H world data, but in good agreement with the  $\mu\text{p}$  value. We motivate an asymmetric fit function, which eliminates line shifts from quantum interference of neighboring atomic resonances.



**Fig. 1 Rydberg constant  $R_\infty$  and proton RMS charge radius  $r_p$ .** Values of  $r_p$  derived from this work (green diamond) and spectroscopy of  $\mu\text{p}$  ( $\mu\text{p}$ : pink bar and violet square) agree. We find a discrepancy of 3.3 and 3.7 combined standard deviations with respect to the H spectroscopy world data (12) (blue bar and blue triangle) and the CODATA 2014 global adjustment of fundamental constants (3) (gray hexagon), respectively. The H world data consist of 15 individual measurements (black circles, optical measurements; black squares, microwave measurements). In addition to H data, the CODATA adjustment includes deuterium data (nine measurements) and elastic electron scattering data. An almost identical plot arises when showing  $R_\infty$  instead of  $r_p$  because of the strong correlation of these two parameters. This is indicated by the  $R_\infty$  axis shown at the bottom.



## Electron Scattering off Proton @ ELPH, Tohoku Univ.

- 1) elastic e+p scattering at lowest-ever  $Q^2$  region
- 2)  $G_E(Q^2)$  at  $0.0003 \leq Q^2 \leq 0.008 \text{ (GeV/c)}^2$
- 3)  $G_E$  is extracted by the **Rosenbluth separation**
- 4) **absolute cross section measurement**  
relative to  $^{12}\text{C}(e,e)^{12}\text{C}$  : sys. err.  $\sim 3 \times 10^{-3}$
- 5)  $E_e = 20 - 60 \text{ MeV}$ ,  $\theta = 30 - 150^\circ$
- 6) constructing of new beam line, and spectrometers
- 7) the experiments will start in 2019