

Development of single Yb^+ and Ba^+ optical clocks and frequency combs towards search for temporal variation of the fine structure constant

Fundamental Physics Using Atoms 2015

December 1, 2015

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(since 2003.9)



Coworkers

- Yb⁺:

Yasutaka Imai (D)

Shunsuke Tomomatsu (M1)

Ren Irie (B4)

S-D3/2, -5/2 clock transition

S-F clock transition

2nd trap

- Ba⁺:

Hiroto Fujisaki (D)

Keisuke Nishida (M2)

Shinya Kawada (B4)

S-D5/2 clock transition

Odd isotope

Clock laser

- Comb:

Masatoshi Mitaki (D)

Solid-state mode-locked laser

Acknowledgement:

Prof. Masao Kitano and Kitano lab.



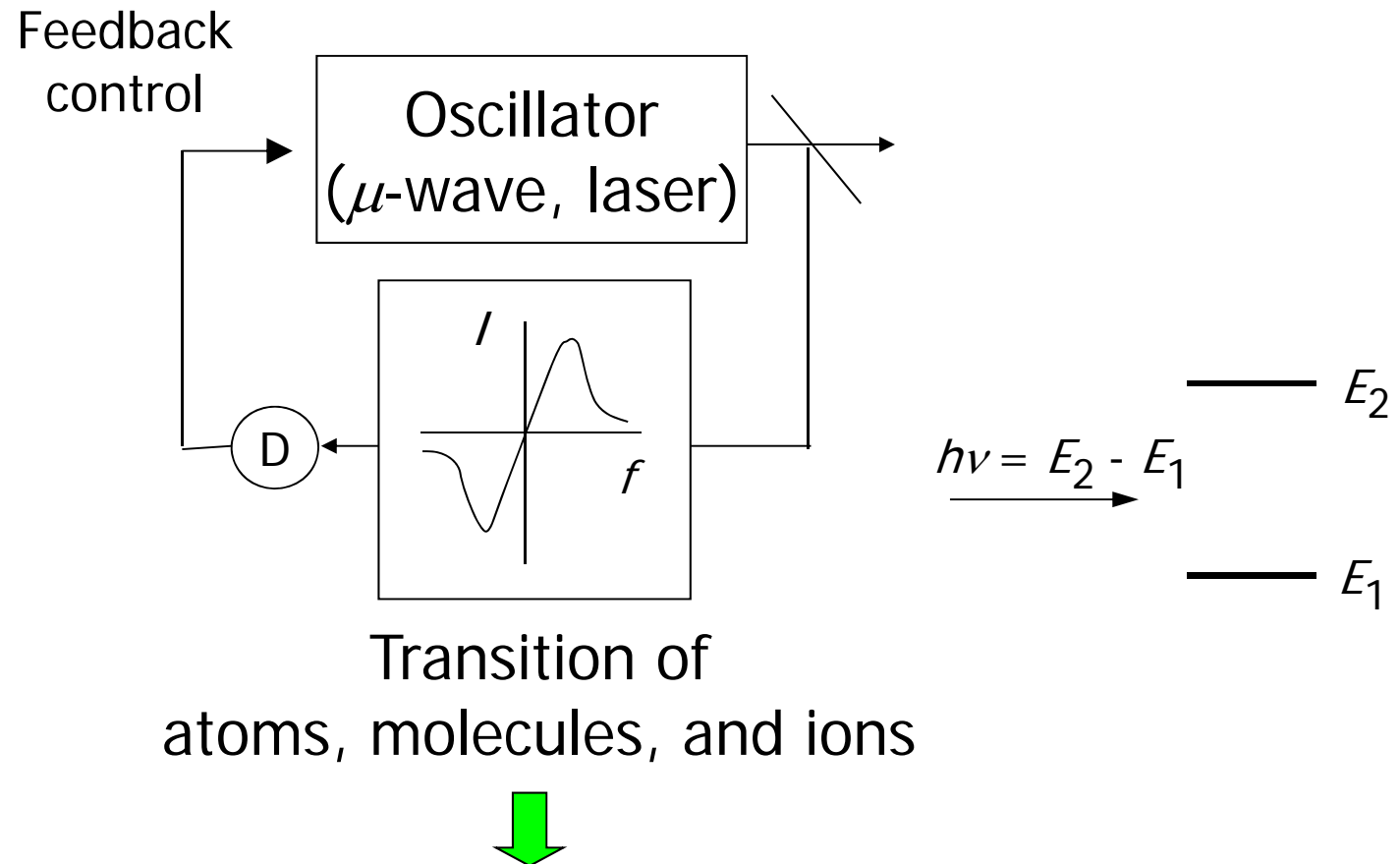
Outline

- Introduction
- Progress after the previous FPUA conference
 - Single-ion spectroscopy of $^2S_{1/2} - ^2D_{3/2}$ in $^{171}\text{Yb}^+$
 - Optical frequency comb
 - Phase locking to laser
 - Diode-laser-pumped Yb:KYW laser

Temporal variation of fundamental constant

- A possibility of temporal variation is predicted in the theories unifying gravity to the other three fundamental interactions.
- Dimensionless constant (independent of units, and temporal variation of standards)
 - α , $\mu = m_p/m_e$
- Observation of the past events
 - Absorption in interstellar matter of rf from quasars
 - null: J. Bagdonaite *et al.*, Science 339, 46 (2013)
 - for: L. K. Webb *et al.*, PRL107, 191101 (2011) (spatial variation?)
 - Observation at one time in the past
 - Model dependent
- Repeatable measurement at present
 - Frequency ratio between atomic clocks of different transitions
 - Requirement: temporal: $<10^{-17}$ /yr $\leftarrow t^{-1}$
(spatial?: + 10^{-20} annual modulation)
Berengut and Flambaum, EPL 97, 20006 (2012)

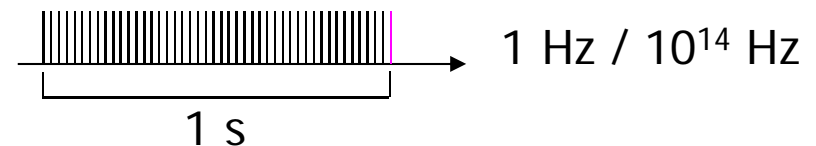
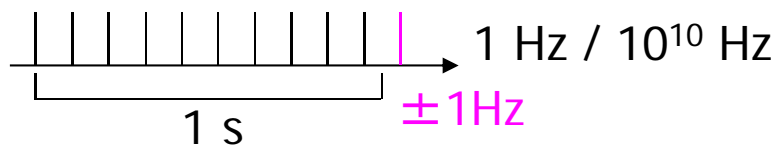
(passive) Atomic clock (frequency standard)



Application to search for temporal variation of α, μ

Advantage of optical clock

- Better uncertainty than that of current standard in microwave
 - Higher the frequency,
better the fractional frequency instability



- Uncertainty: $<10^{-17}$
 - single-ion clock $\text{Al}^+ : 7.0 \times 10^{-18}$ Chou *et al.*, PRL 104, 070802 (2010)
 $\text{Yb}^+ : 3 \times 10^{-18}(?)$ PTB G., ICOLS 2015
 - optical lattice clock $\text{Sr} : 7 \times 10^{-18}$ Ushijima *et al.*, Nat. Photon. 9, 185 (2015)

Advantage of optical clock

- Model independent, concerning α variation

Electronic transition frequency f_E $f_E = \text{const} \cdot (R_\infty c) \cdot F(\alpha)$

Temporal variation of f_E $\frac{1}{f_E} \frac{df_E}{dt} = \frac{1}{R_\infty c} \frac{d(R_\infty c)}{dt} + K \frac{1}{\alpha} \frac{d\alpha}{dt}, \quad K \equiv \frac{\alpha}{F(\alpha)} \frac{dF(\alpha)}{d\alpha}$

Frequency ratio f_{E1}/f_{E2}
$$\boxed{\frac{1}{f_{E1}/f_{E2}} \frac{d(f_{E1}/f_{E2})}{dt} = (K_{E1} - K_{E2}) \frac{1}{\alpha} \frac{d\alpha}{dt}}$$

$F(\alpha)$: Relativistic correction

sensitivity to α depends on the transition

Rydberg-constant ($R_\infty c$ in Hz) variation is the same in all transitions

Sensitivity K to the temporal variation of α

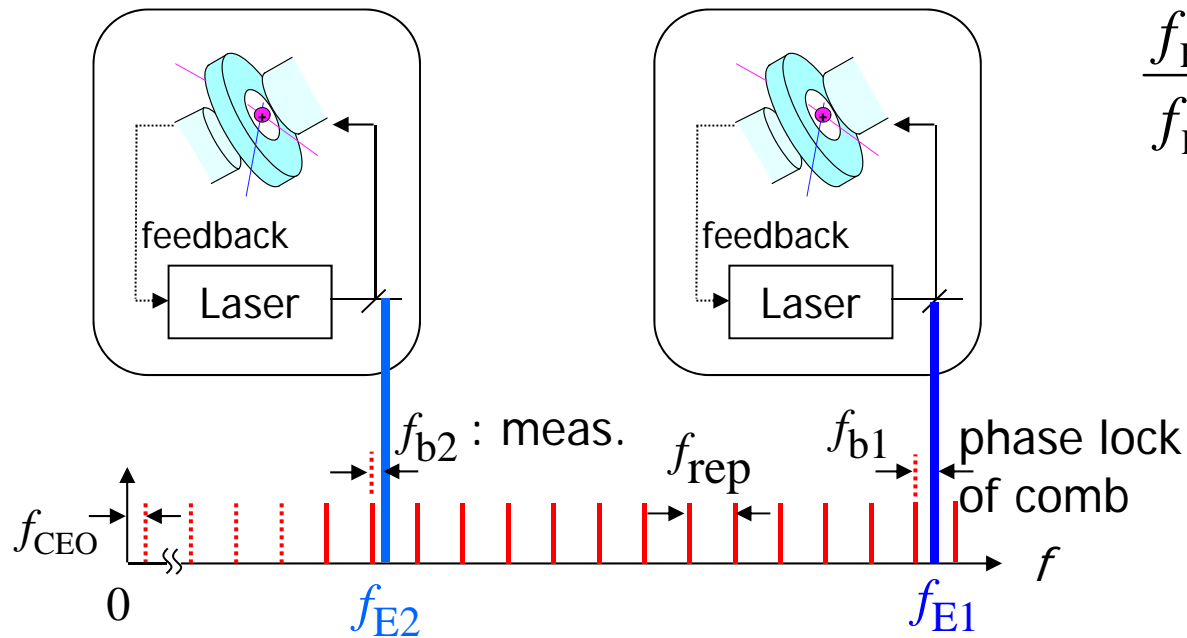
- K : Sensitivity to variation of α
 - Relativistic effect
 - Large interval between fine structure : high sensitivity
 - Scale with $(Z\alpha)^2$
 - Atoms of heavier mass

atom	transition	A
H	$1s_{1/2}-2s_{1/2}$	2.4×10^{-5}
Ca	$1S_0-^3P_0$	0.016
Sr	$1S_0-^3P_0$	0.06
Yb	$1S_0-^3P_0$	0.31
Hg	$1S_0-^3P_0$	0.81

ion	transition	K
Al ⁺	$1S_0-^3P_0$	0.008
In ⁺	$1S_0-^3P_0$	0.18
Ca ⁺	$2S_{1/2}-^2D_{5/2}$	0.15
Sr ⁺	$2S_{1/2}-^2D_{5/2}$	0.43
Ba ⁺	$2S_{1/2}-^2D_{3/2}$	2.5
	$2S_{1/2}-^2D_{5/2}$	2.4
Yb ⁺	$2S_{1/2}-^2D_{3/2}$	1.00
	$2S_{1/2}-^2D_{5/2}$	1.03
	$2S_{1/2}-^2F_{7/2}$	-6.0
Hg ⁺	$2S_{1/2}-^2D_{3/2}$	-1.5
	$2S_{1/2}-^2D_{5/2}$	-2.9

V. Flambaum and V. Dzuba, Can J. Phys. 87 25 (2009)

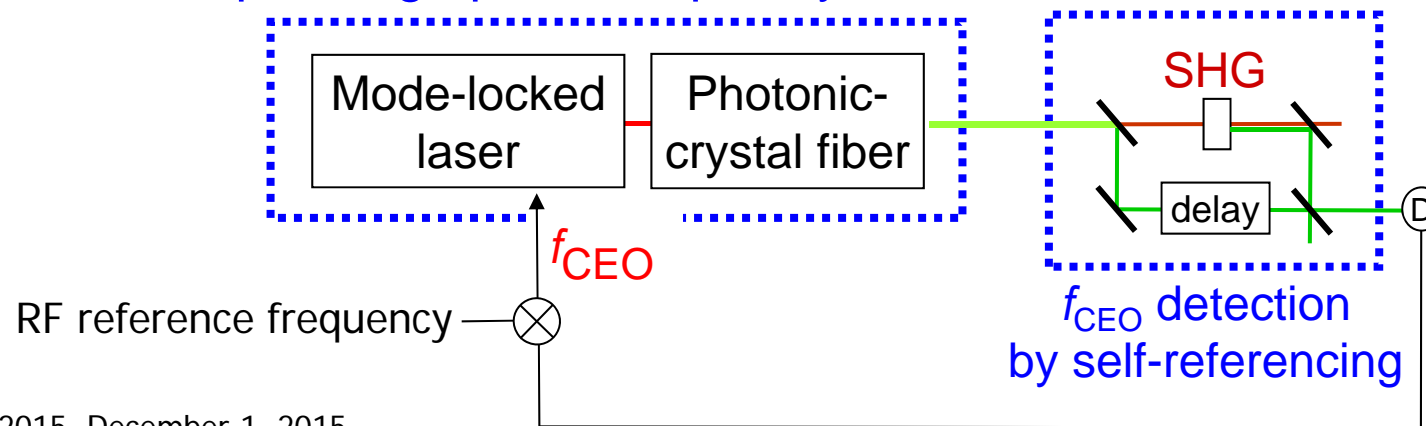
Optical frequency ratio measurement



$$\frac{f_{\text{E2}}}{f_{\text{E1}}} = \frac{n_2 f_{\text{rep}} + f_{\text{CEO}} + f_{\text{b2}}}{n_1 f_{\text{rep}} + f_{\text{CEO}} + f_{\text{b1}}}$$

- Value of RF reference frequency is not required.
- Stability of 10^{-12} required in RF reference for uncertainty 10^{-18} ($n \sim 10^6$)
- Uniformity of comb spacing is proved in 10^{-19} .

Octave-spanning optical frequency comb



Constraints from comparisons of atomic clocks

$$\frac{1}{\alpha} \frac{d\alpha}{dt}$$

$^{27}\text{Al}^+ / ^{199}\text{Hg}^+$

$-1.6(2.3) \times 10^{-17} \text{ /yr}$ at NIST

T. Rosenband *et al.*,

Science, 319, 1808 (2008).

Please refer to fig. 4 in

R.M. Godun *et al.*, PRL 113, 210801 (2014)

incl. $^{171}\text{Yb}^+ \text{E2} / ^{171}\text{Yb}^+ \text{E3}$

$-0.7(2.1) \times 10^{-17} \text{ /yr}$ at NPL

$-2.0(2.0) \times 10^{-17} \text{ /yr}$ at PTB

$$\frac{1}{\mu} \frac{d\mu}{dt}$$

Comparison to Cs hyperfine

incl. $^{171}\text{Yb}^+ \text{E2} / ^{171}\text{Yb}^+ \text{E3}$

$0.2(1.1) \times 10^{-16} \text{ /yr}$ at NPL

$-0.5(1.6) \times 10^{-16} \text{ /yr}$ at PTB

Please refer to fig. 3 in

N. Huntemann *et al.*, PRL 113, 210802 (2014)

Advantage of Yb⁺ and Ba⁺ on search for temporal variation of α

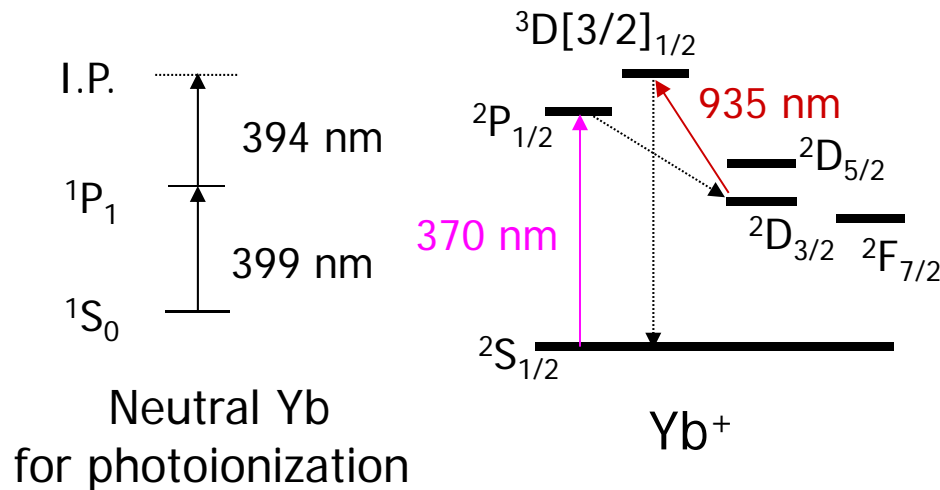
- Yb⁺:
 - Measurement on one ion in one trap
 - Close investigation of perturbations
 - Cancellation of gravitational effect
gravitational red shift: $10^{-18}/\text{cm}$
 - S-F: sensitivity twice as large as S-D_{5/2} in Hg⁺
 - S-D_{3/2,5/2}: variation vanished
 - Comparison with other labs (PTB, NPL)
- Ba⁺
 - Odd isotope: insensitive to electric quadrupole shift
 - Low uncertainty
 - Increase of number

- Frequency ratio measurement on three transitions in Yb⁺ & Ba⁺

- Ratio measurement among three
 - Absolute variation
 - Extract temporal variation of Rydberg constant

	$K_1 - K_2$
$\text{Yb}^+ (^2S_{1/2} - ^2D_{3/2}) - \text{Yb}^+ (^2S_{1/2} - ^2D_{5/2})$	-0.03
$\text{Yb}^+ (^2S_{1/2} - ^2D_{3/2}) - \text{Yb}^+ (^2S_{1/2} - ^2F_{7/2})$	7.1
$\text{Ba}^+ (^2S_{1/2} - ^2D_{3/2}) - \text{Yb}^+ (^2S_{1/2} - ^2F_{7/2})$	8.5
$\text{Ba}^+ (^2S_{1/2} - ^2D_{3/2}) - \text{Yb}^+ (^2S_{1/2} - ^2D_{3/2})$	1.5

Yb⁺



Clock transitions:

411 nm:	$2S_{1/2} - 2D_{5/2}$	$\tau = 7 \text{ ms}$
435 nm:	$2S_{1/2} - 2D_{3/2}$	$\tau = 52 \text{ ms}$
467 nm:	$2S_{1/2} - 2F_{7/2}$	$\tau \sim 4000 \text{ d}$

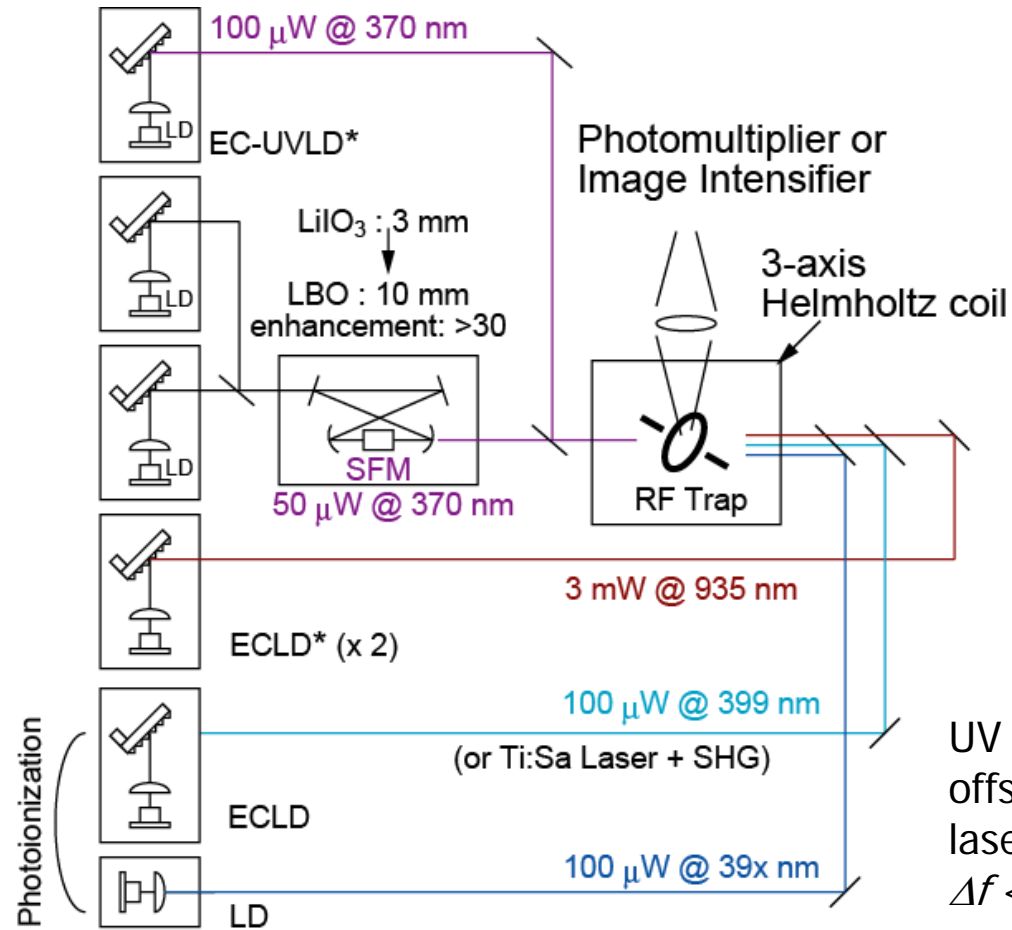
Roberts et al PRA60, 2867 (1999)

Tamm et al, PRA80, 043403(2006)

Huntemann et al., PRL 108, 090801(2012)

- Isotope 171 ($I=1/2$)
 - $m_F=0 \rightarrow m_F'=0$, no 1st-order Zeeman shift
 - simplest hyperfine structure because of $I=1/2$
 - (only $^{199}\text{Hg}^+$ has $I=1/2$)
- Small system with simple light source
 - Secondary frequency standard
PTB, NPL
 - Quantum information processing
U. Maryland, etc.

Setup for laser cooling of Yb⁺



$$2r_0 = 0.8 \text{ mm}$$

$$\Omega/2\pi = 13.5 \text{ MHz}$$

UV SF radiation is frequency stabilized by offset locking of one of the fundamental laser to the clock laser at 822 nm.
 $\Delta f < \pm 200 \text{ kHz @ } 370 \text{ nm}$

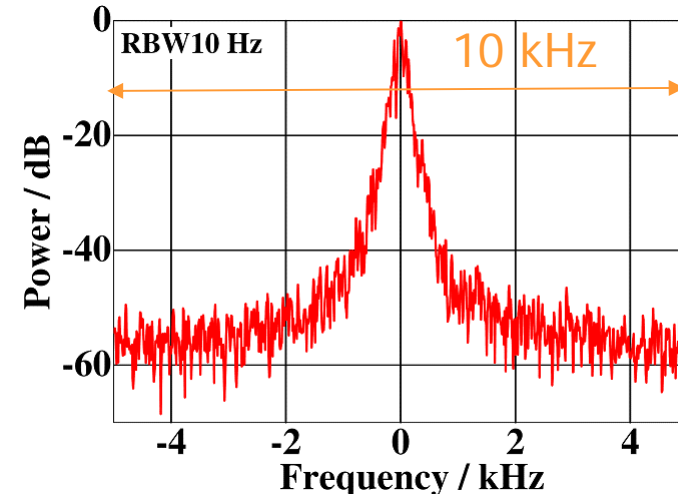
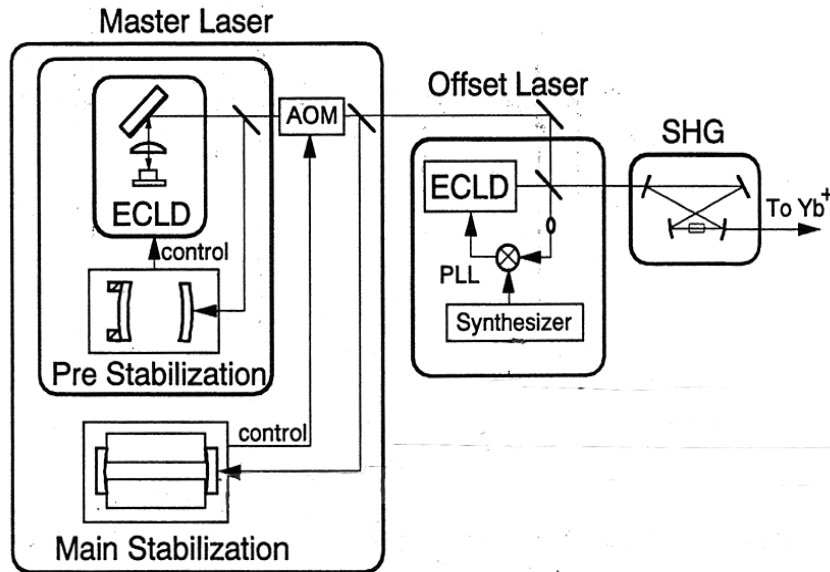
ECLD: Extended-cavity laser diode
 * Stabilized with temperature controlled cavity

We acknowledge Prof. Kawakami and Nichia Ltd.
 for providing us UV-LDs.

FPUA2015, December 1, 2015

SFM: Sugiyama et al., Applied Optics, 49, 5510 (2010)
 Photoionization: Onoda et al., Applied Physics B (2011)

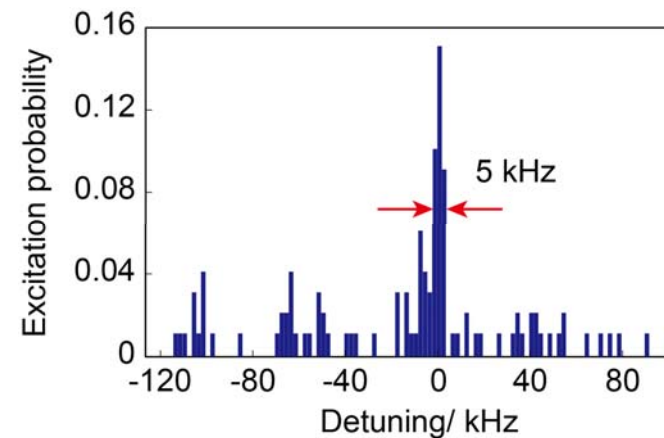
Spectroscopy of the $^2S_{1/2} - ^2D_{5/2}$ clock transition at 411 nm in $^{174}\text{Yb}^+$



Beatnote of two independent systems
Linewidth < 300 Hz

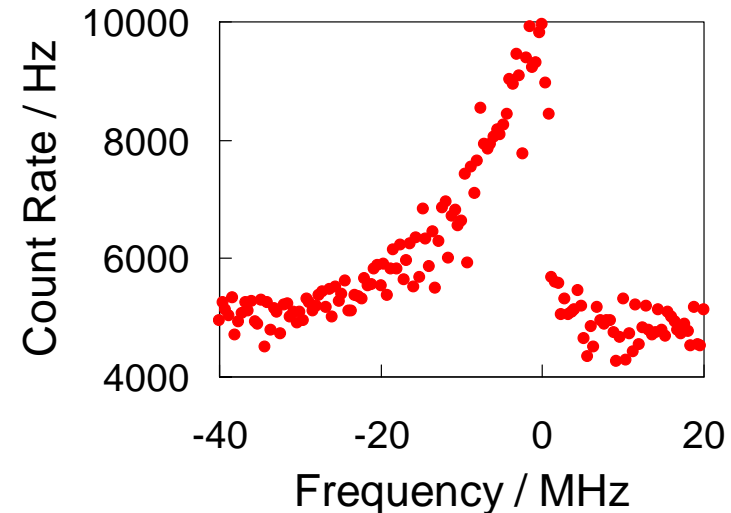
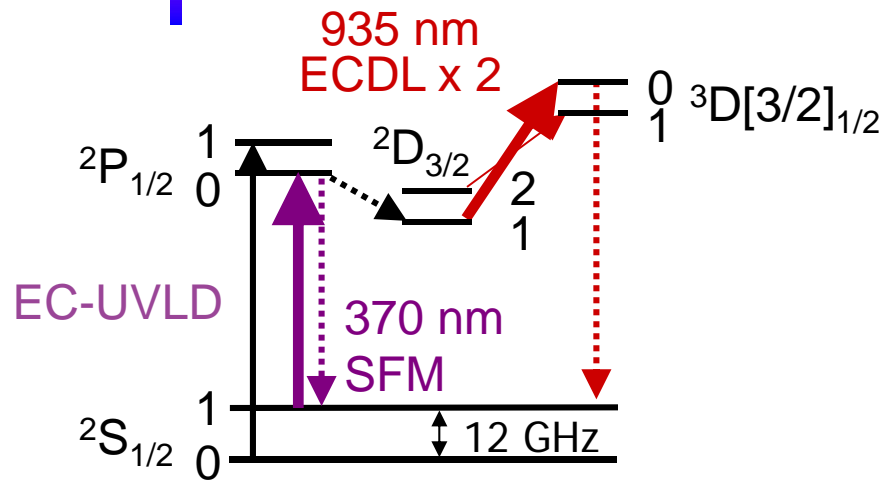
Pre: $\Delta_{\text{cavity}} = 15$ MHz (super inver)
Transparent fringe side
fast: current,
slow: PZT

Main: $\Delta_{\text{cavity}} = \sim 80$ kHz (ULE)
FM sideband
fast: AOM
slow: PZT (Pre stab. cavity)



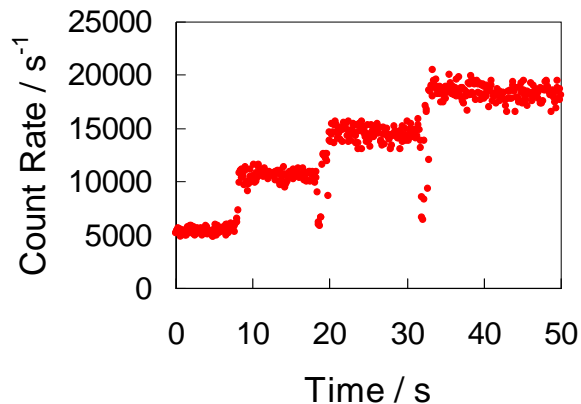
$m_j = +1/2 \rightarrow +5/2$
Step: 2 kHz

Laser cooling of single $^{171}\text{Yb}^+$



Spectrum of cooling transition

Temperature: ~ 13 mK



One-by-one loading
with photoionization

~ 6000 cps/ion

- Micromotion minimization by elimination of stray DC field

← rf-photon cross correlation detection

Berkeland et al., JAP 83, 5025 (1998)

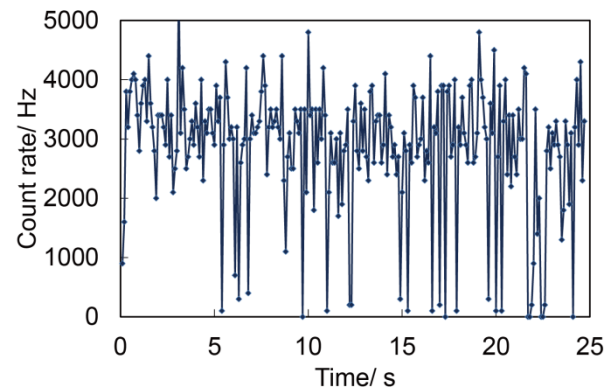
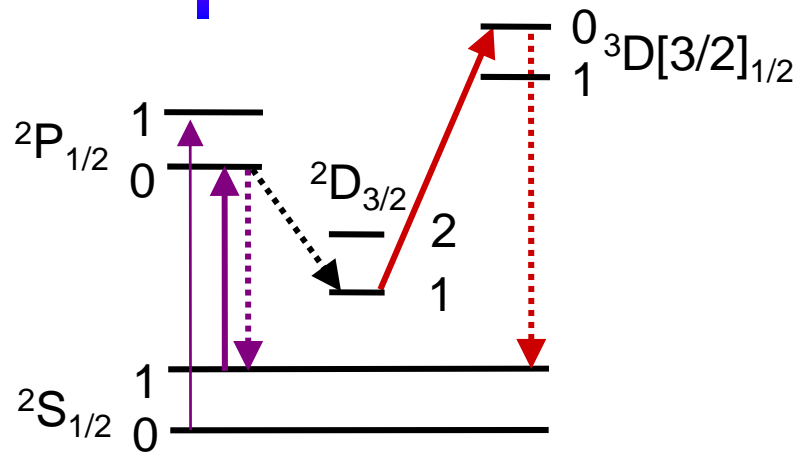
- Elimination of dark states

← Appropriate magnetic field and laser polarization

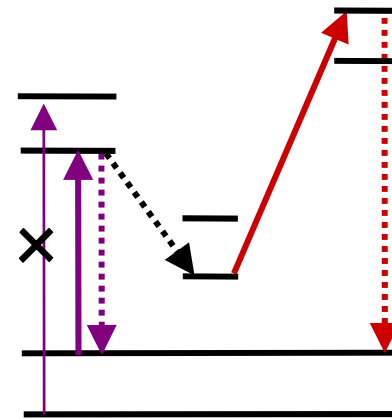
Berkeland et al., PRA 65, 033413 (2002) 15

Driving of the $^2S_{1/2} - ^2D_{3/2}$ clock transition at 435 nm

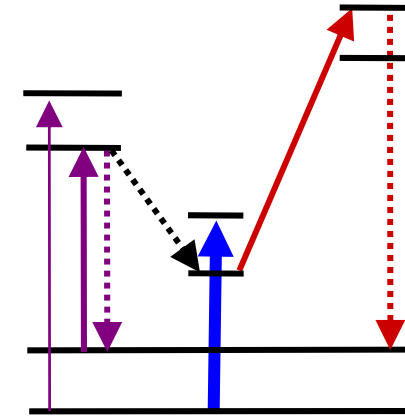
Imai



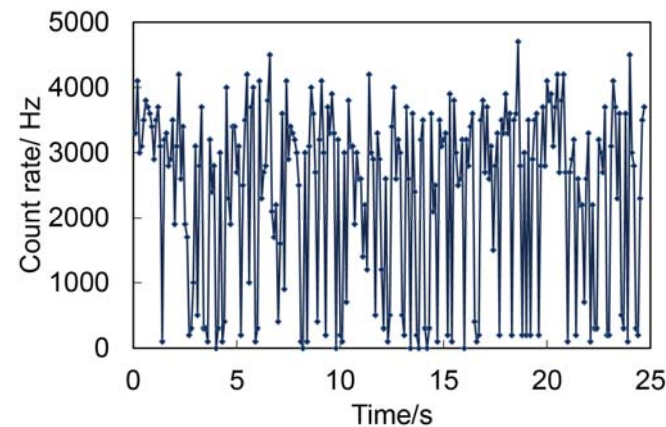
without sub-repumping
($D_{3/2} F=2 \rightarrow [3/2]_{1/2} F=1$)
spontaneous decay
from $P_{1/2} F=1$ to $D_{3/2} F=2$



Initialization to $S_{1/2} F=0$



Clock irradiation
/Detection



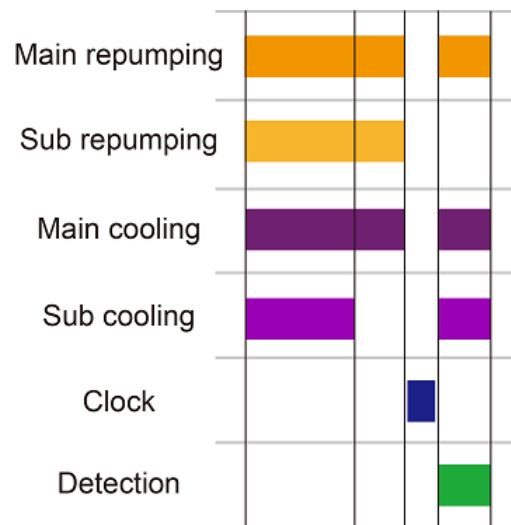
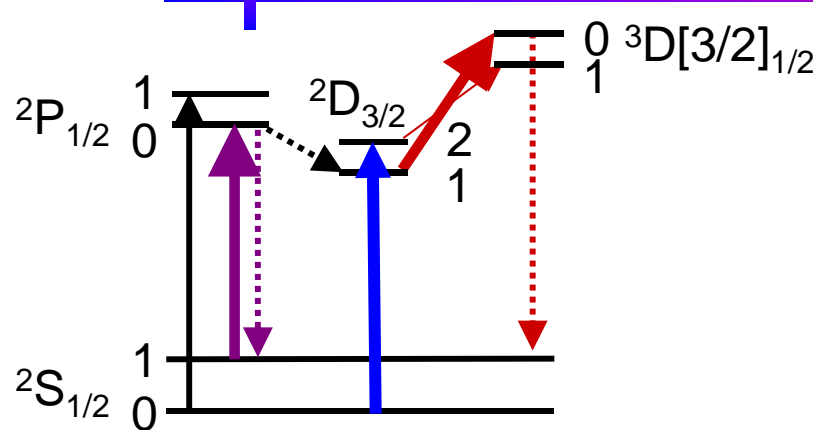
Quantum jumps with clock transition

Clock : 10 ms
Detection: 10 ms

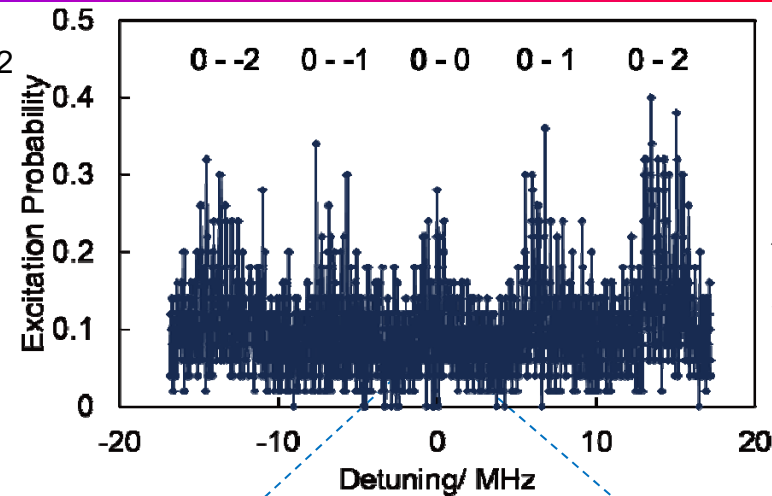
Ref. D. Engelke and C. Tamm, Europhys. Lett. 33, 347 (1996)

V. Buehner and C. Tamm, Phys. Rev A 61, 061801(R)(2000)

Spectra of $^2S_{1/2}$ (F=0) – $^2D_{3/2}$ (F=2) transition in single $^{171}\text{Yb}^+$

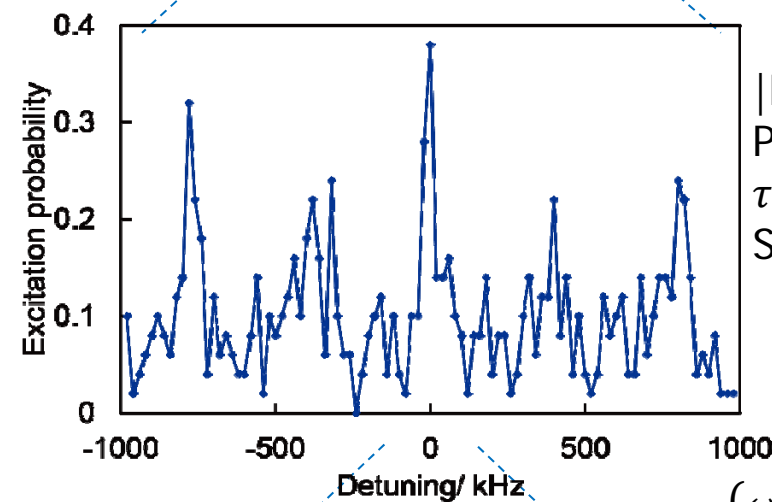


Initialization : 10 ms
 Clock irradiation: 10 ms
 Detection: 10 ms
 Averaging 50 trials



$|B| \sim 830 \mu\text{T}$
 $P_{\text{clock}} = 60 \mu\text{W}$
 $\tau = 10 \text{ ms}$
 Step: 20 kHz

All Zeeman components

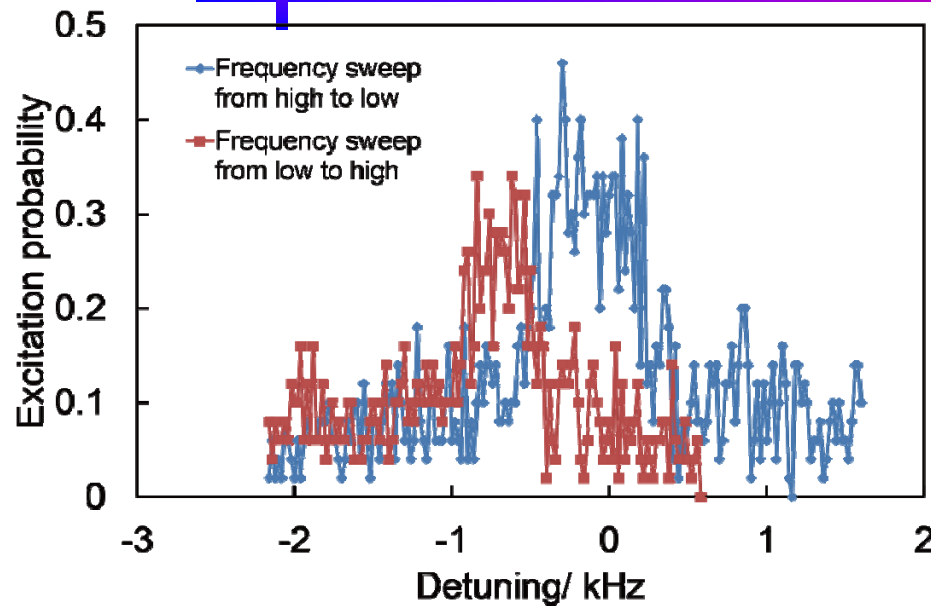


$|B| \sim 830 \mu\text{T}$
 $P_{\text{clock}} = 60 \mu\text{W}$
 $\tau = 10 \text{ ms}$
 Step: 20 kHz

mF=0 – 0 transition with motional sidebands

$\begin{cases} \omega_{\text{rad}} \cong 370 \text{ kHz} \\ \omega_{\text{ax}} \cong 760 \text{ kHz} \end{cases}$

High-resolution scanning of carrier spectrum



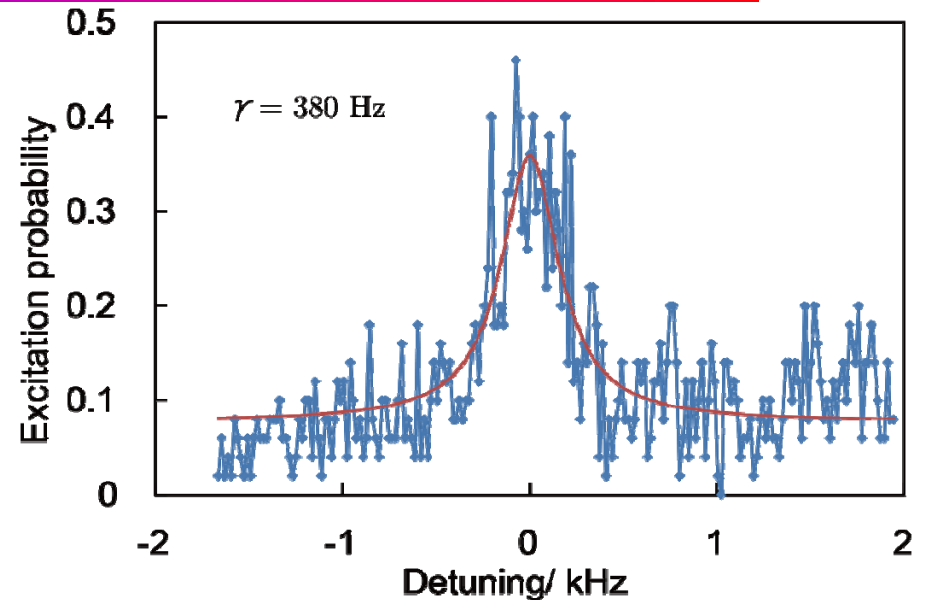
$m_F = 0 - m_{F'} = 0$ carrier spectrum

Feed forward compensation of drift of clock laser: +32 Hz/s

Sweep direction:

Red: low to high

Blue: high to low



After further adjustment to agree the spectra obtained by two sweep directions

FWHM ~ 380 Hz

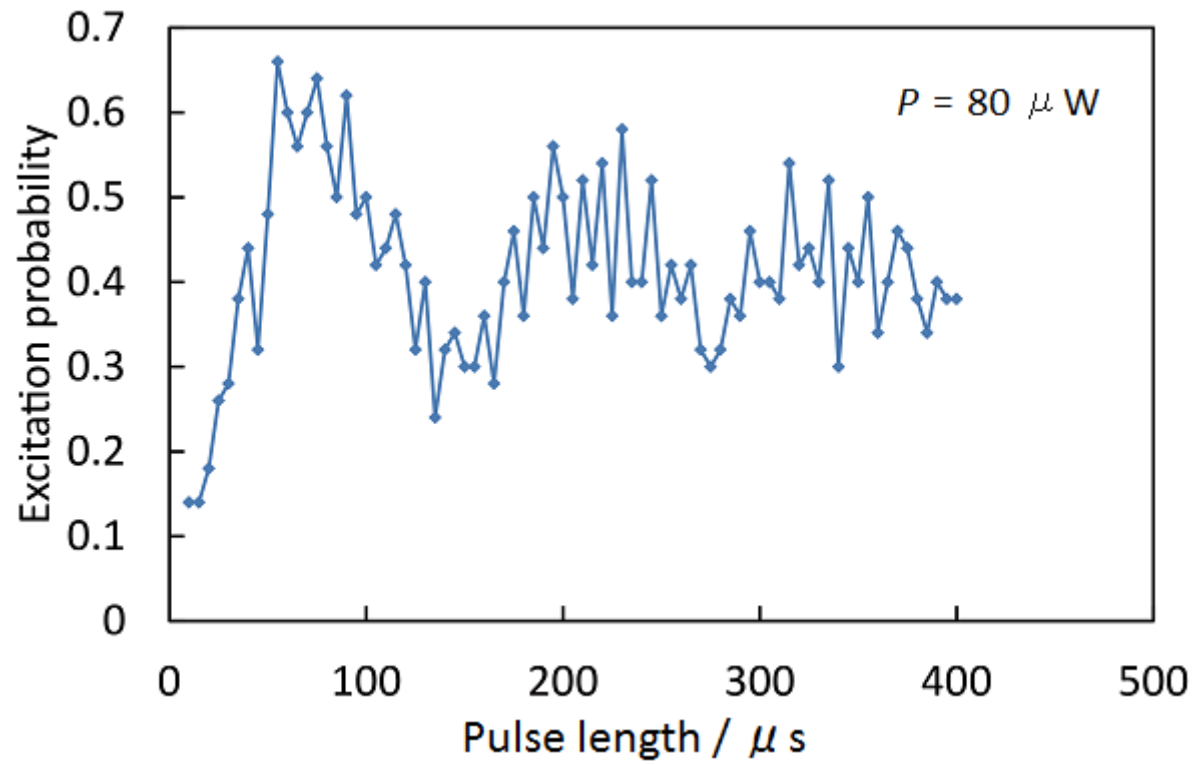
limited by laser linewidth

$|B| \sim 830 \mu\text{T}$, $P_{\text{clock}} = 10 \mu\text{W}$

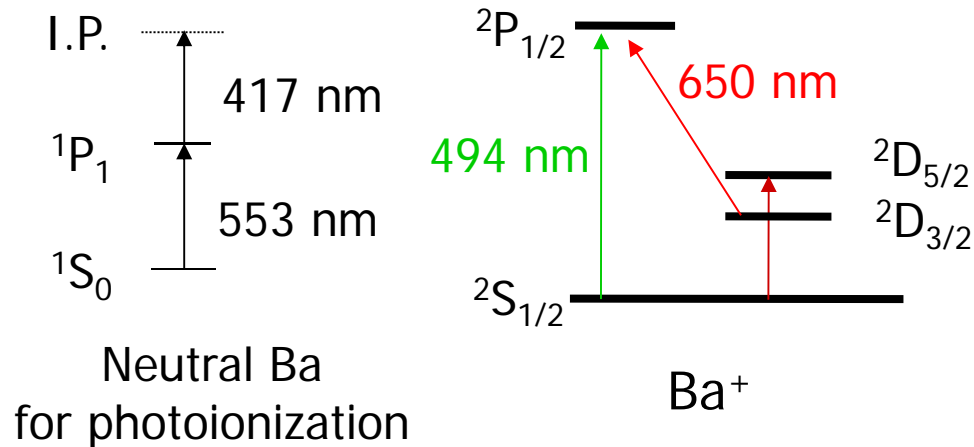
$\tau = 10$ ms, Step: 20 Hz

18

Finding π pulse



Ba⁺



● Clock transitions

1.762 μm : $2S_{1/2} - 2D_{5/2}$ $\tau = 30$ s

2.052 μm : $2S_{1/2} - 2D_{3/2}$ $\tau = 17$ s

ECLD, DFM
DFM

● Isotope 135, 137 ($I=3/2$)

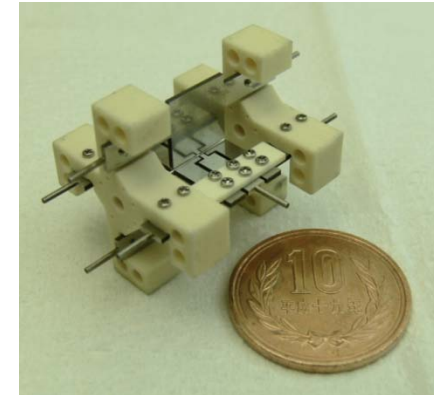
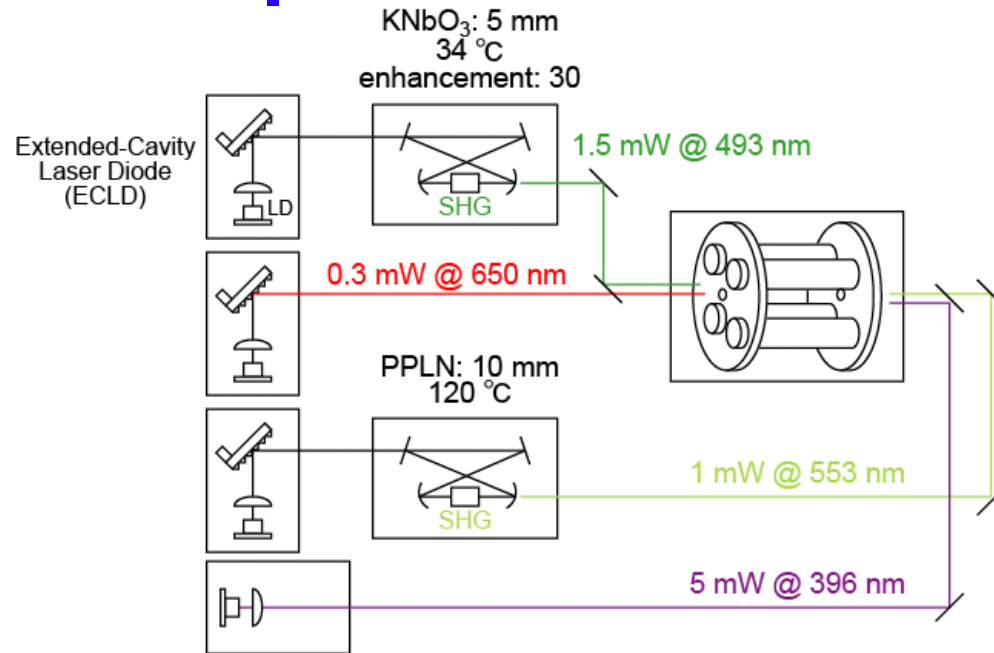
- $2D_{3/2}$, $F=0$: very small electric quadrupole moment (largest factor of uncertainty in ions with alkali-metal like level scheme)
- $m_F=0 \rightarrow m_F'=0$, no 1st-order Zeeman shift

Blinov Gr. & Fortson Gr. (U. Washington)

Blatt Gr. (U. Insbrueck)

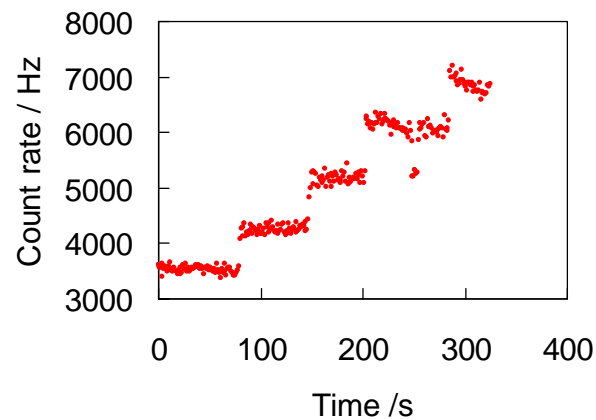
U. Groningen, U. Singapole

Setup for laser cooling of Ba⁺



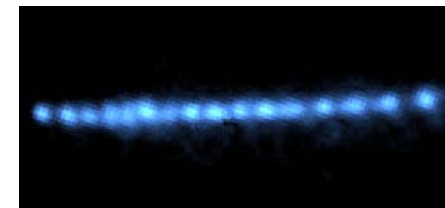
$$2r_0 = 0.8 \text{ mm}, 2z_0 = 1.0 \text{ mm}$$

Segmented-electrode linear RF trap
(acknowledgement: Prof. Urabe (U.Osaka),
Dr. Hayasaka (NICT))



One-by-one loading of ¹³⁸Ba⁺

with photoionization



String of 16 ¹³⁸Ba⁺ ions

Single $^{138}\text{Ba}^+$ spectroscopy of the $^2\text{S}_{1/2} - ^2\text{D}_{5/2}$ transition at 1762 nm

Fujisaki

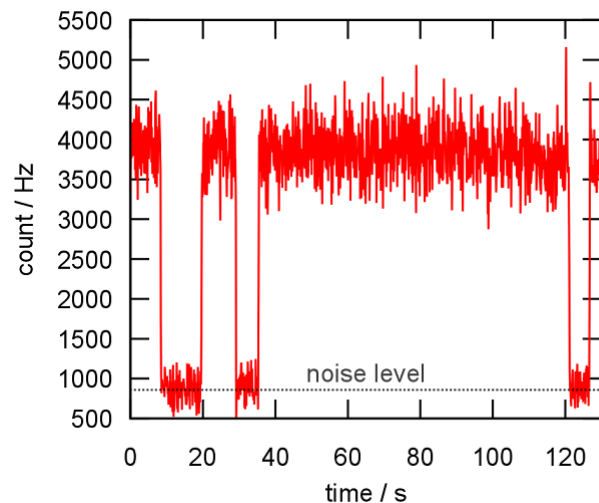
● Clock laser

1) ECLD at 881 nm locked to a high-finesse cavity

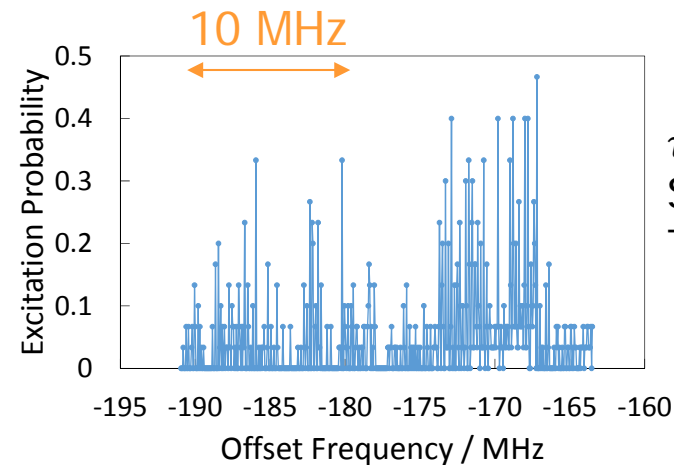
relative linewidth ~ 10 Hz

comparison with comb < 200 Hz

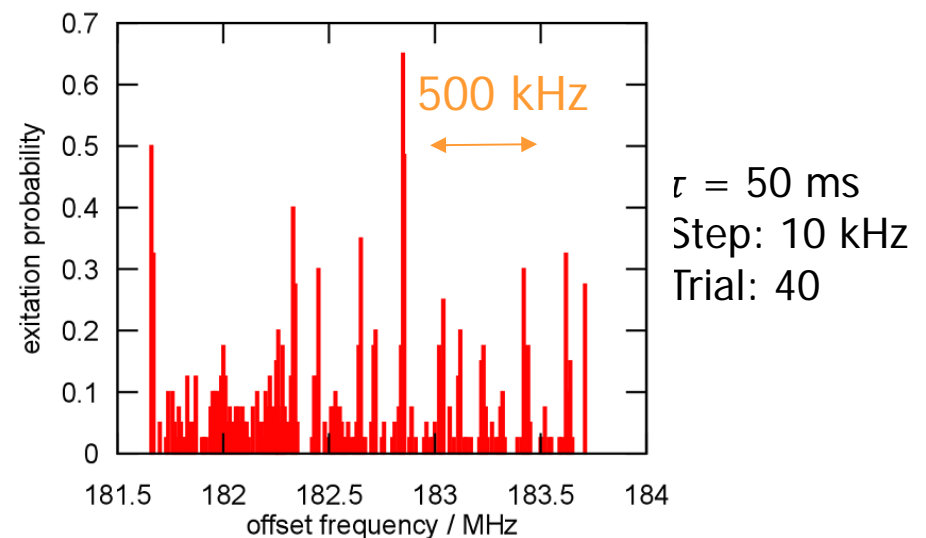
2) Phase locking of ECLD at 1762 nm using second-harmonics.



Quantum jumps of single $^{138}\text{Ba}^+$ by using linewidth-narrowed 1762-nm ECLD

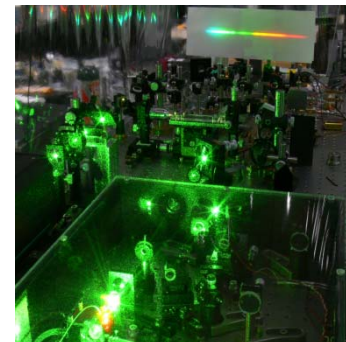
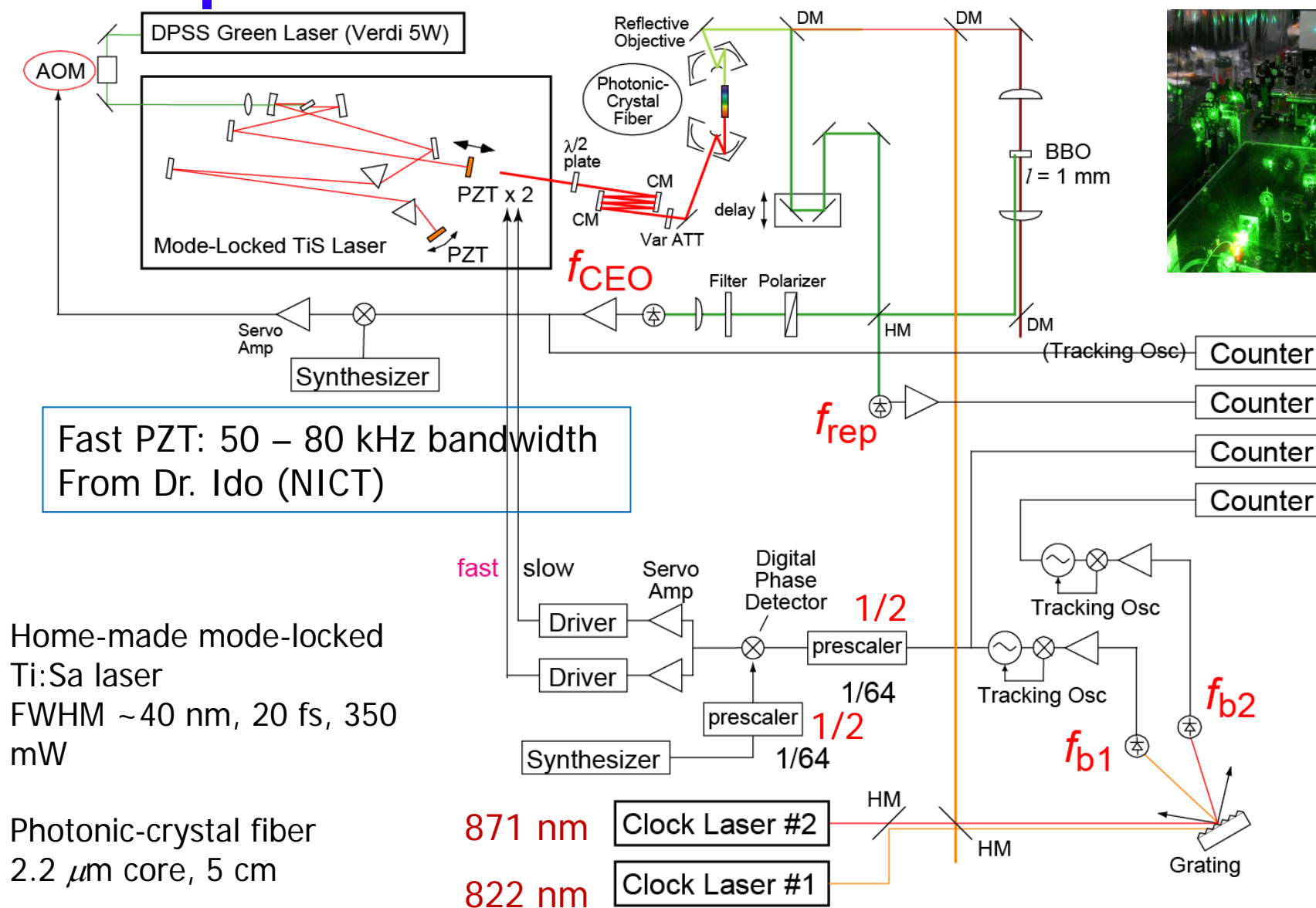


Zeeman components?



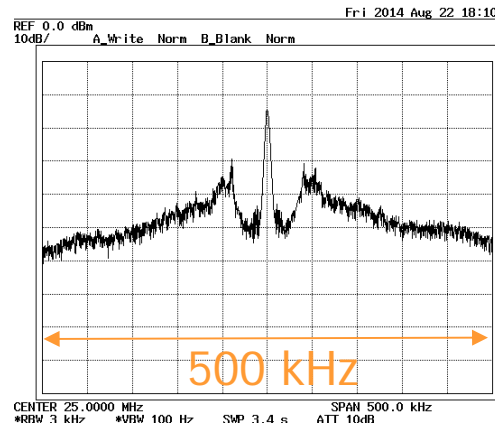
Motional sideband

Setup for frequency ratio measurement using Ti:Sa laser

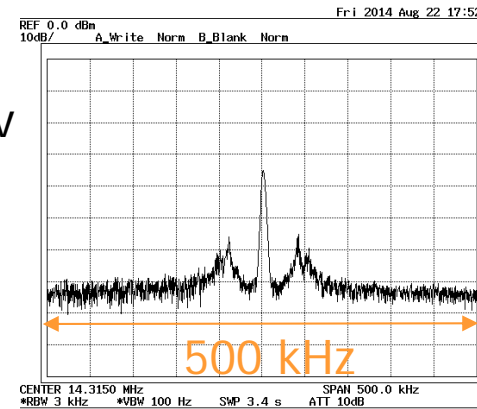


Phase locking of comb to laser

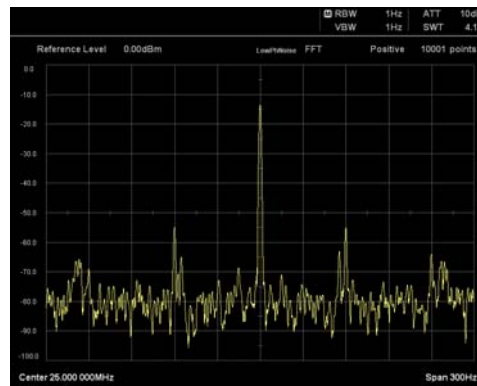
Hatanaka



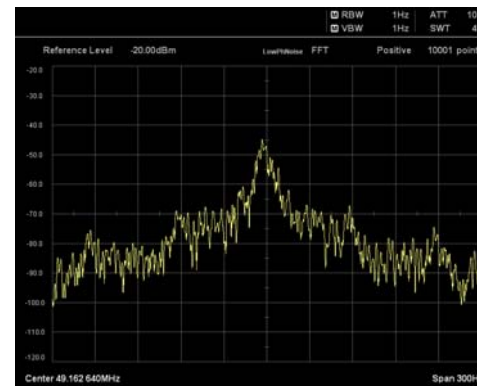
y: 10 dB/div



Span: 500 kHz, RBW: 30 Hz



Span: 300 Hz, RBW: 1 Hz



FWHM:
<10 Hz

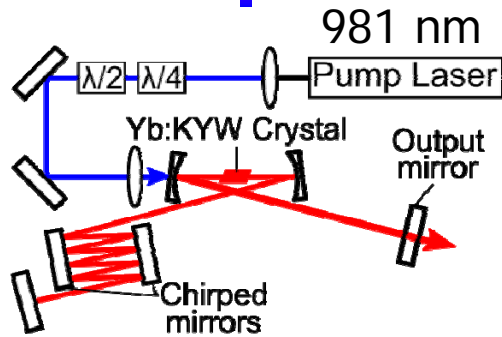
822 nm : Phase locked 871 nm : measured

Power concentration to carrier: 92 %

Residual phase noise: 0.32 rad (10 Hz - 1 MHz)

Laser-diode pumped Kerr-lens mode-locked Yb:KYW laser

Mitaki



Goal:

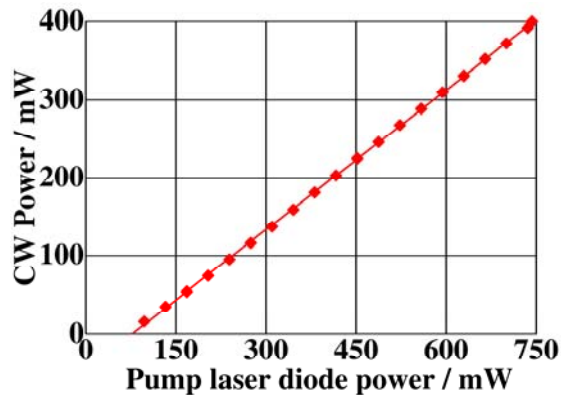
- Long-term continuous operation with low running cost
- Alternative to fiber laser

Advantage:

- Low heat generation

Comb: Meyer *et al.*, Appl. Phys. B 112, 5765 (2013).

High freq: Endo *et al.*, Opt. Express, 20, 12191 (2012).



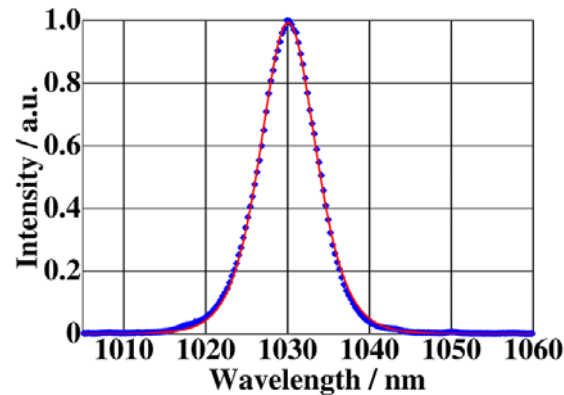
CW oscillation

Max. power

(output coupler: 90%):

400 mW @ 750 mW pump

Slope efficiency: 59 %



Mode-locked oscillation

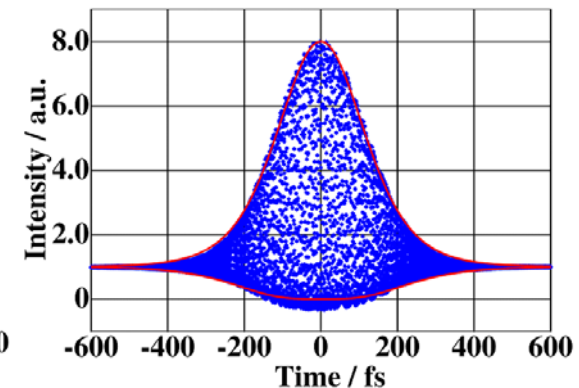
Output coupler 92.5 %, freq ~90 MHz

Spectral width:

8.0 nm @ 360 mW

Pulse duration:

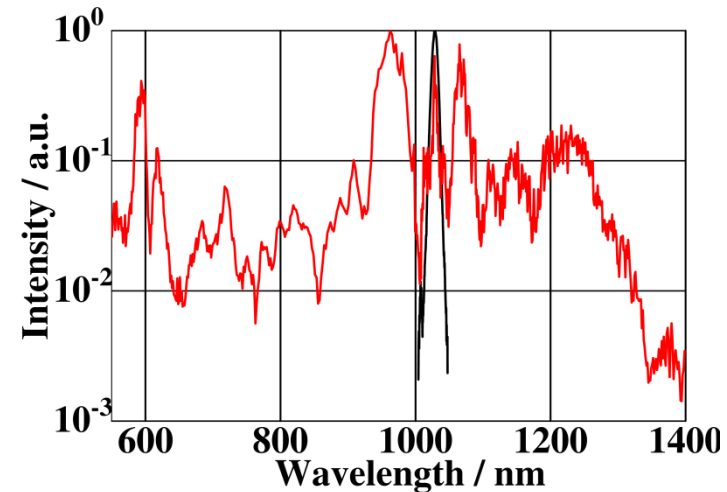
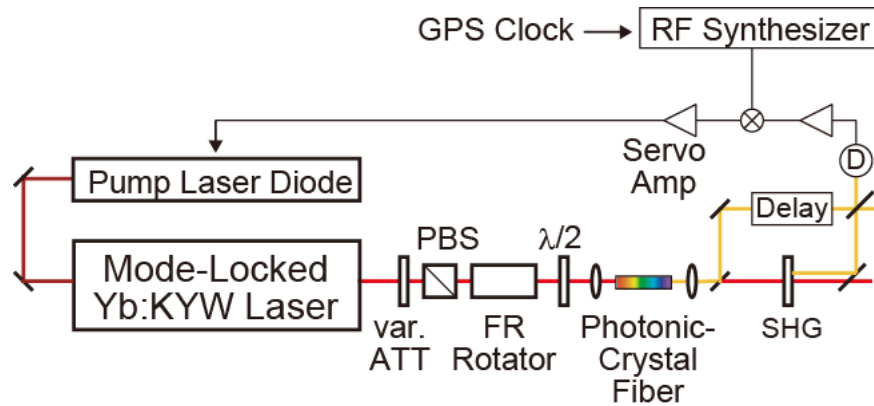
165 fs @ 360 mW



Continuous mode-locking ~1 month

with temperature-controlled base plate

Detection of f_{CEO} with self-referencing technique



Self-referencing technique

Second-harmonic of low-frequency wing

$$2f(n) = 2 \times (n f_{\text{rep}}) + 2f_{\text{CEO}}$$

high-frequency wing of fundamental

$$f(2n) = 2n f_{\text{rep}} + f_{\text{CEO}}$$

→ beat frequency = f_{CEO}

Jones *et al.*, Science 288, 635 (2000)

Spectral broadening with
photonic-crystal fiber

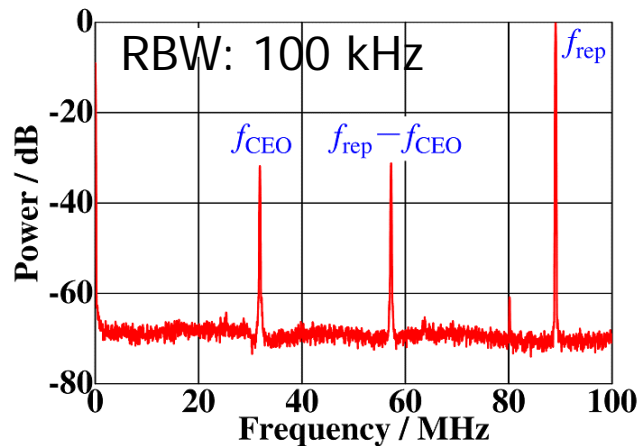
Core: 3.3 μm , Length: 5 cm,

Input: 360 mW

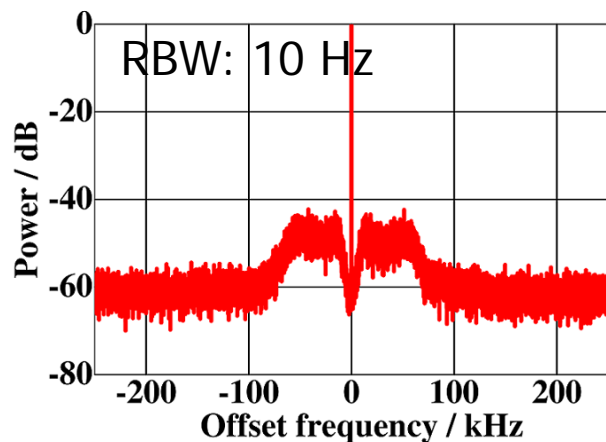
Output: 200 mW

Fiber amplifier not required

Phase locking of f_{CEO}



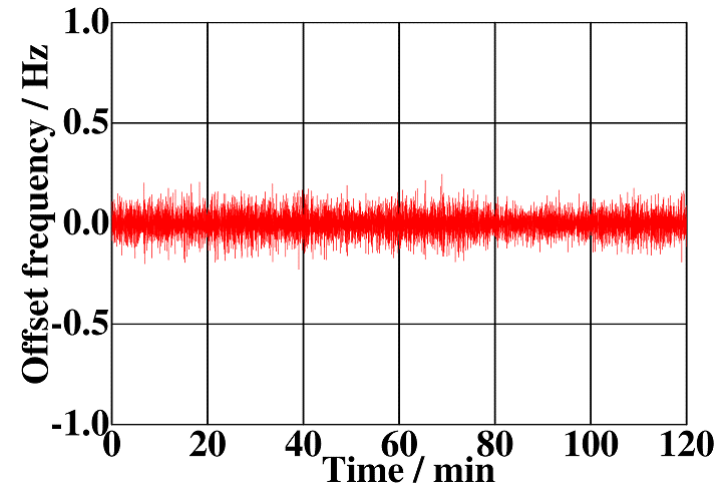
Free running



phase-locked

Beat at f_{CEO}

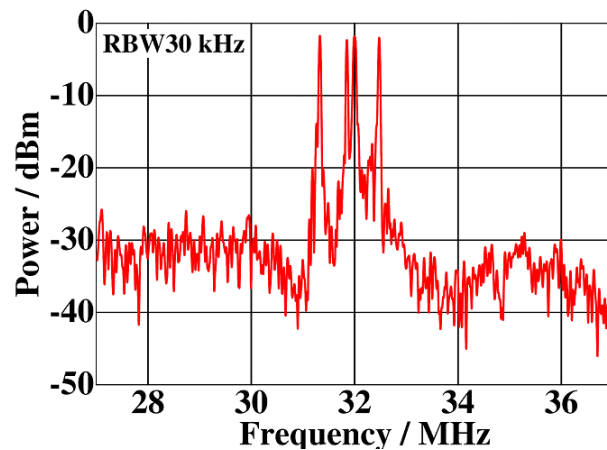
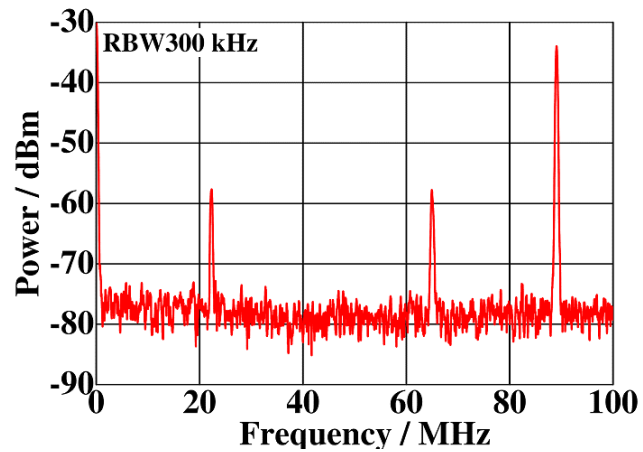
- Residual phase noise:
0.48 rad (1 Hz—1 MHz)
- Concentration in carrier:
90 % in 25-Hz BW
within 250-kHz span



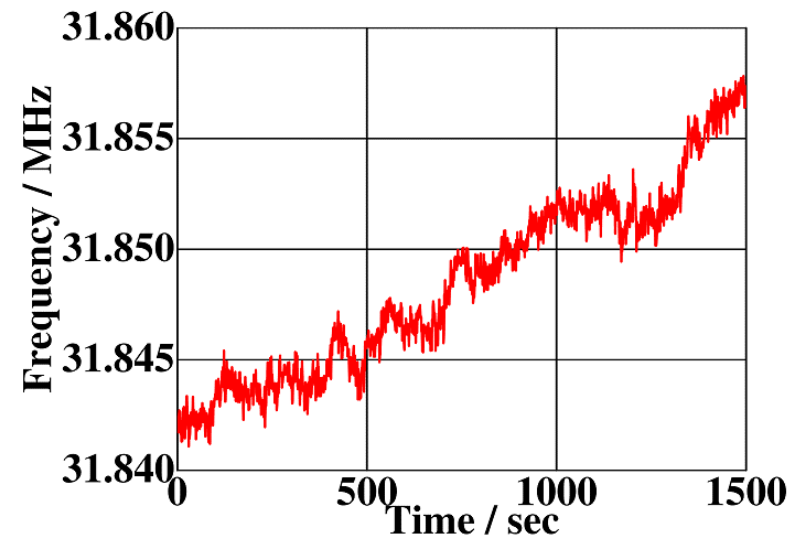
Frequency counting of f_{CEO}

Continuous locking: 2 hours
Gate: 1 s

Frequency measurement



Beat with
a clock laser at 871 nm



Frequency counting of fbeat

Gate: 1 s

Drift: 10 Hz/s, Fluctuation: 3×10^{-12}



Summary

- Development of single Yb⁺ and Ba⁺ ion optical clock, and frequency comb
 - Application: search for temporal variation of the fine structure constant

Current status

- Single-ion spectroscopy: ¹⁷⁴Yb⁺ S-D_{5/2} and ¹⁷¹Yb⁺ S-D_{3/2} transition
 - Resolution: 380 Hz, laser-linewidth limit (<300 Hz)
- Ba⁺: single-ion spectroscopy of 1.76 μm S-D_{5/2} transition
- Comb: Ti:Sa: phase locking to laser completed
Yb:KYW laser : frequency measurement system

Next step

- ★ Single ¹⁷¹Yb⁺ spectroscopy
 - Improvement of linewidth and stability of clock lasers → higher resolution
 - S-F transition
- ★ Second system → Comparison between same or different transitions
 - Ba⁺: single-ion spectroscopy improved, Better cooling of odd isotopes
 - Comb: Yb:KYW: phase locking to laser